# STATIONARY MEASURES ON THE CIRCLE FROM HYPERBOLIC SURFACES WITH CUSPS CANNOT BE STRAIGHTENED BY QUASI-SYMMETRIES 

AITOR AZEMAR AND VAIBHAV GADRE


#### Abstract

Stationary measures on the circle that arise from a large class of random walks on the fundamental group of a finite-area complete hyperbolic surface with cusps are singular with respect to the Lebesgue measure. In particular, it is sufficient for singularity that a stationary measure satisfies an exponential decay for cusp excursions with excursions being measured in the path metric on horocycles bounding cusps. In this note, we settle a conjecture of McMullen by proving that the singularity of stationary measures satisfying such exponential decay is quasi-symmetrically stable, that is under push-forward by any quasi-symmetry of the circle the measure remains singular.


## 1. Introduction

Two intervals on the circle are said to be adjacent and equal if they share an endpoint and have the same length. A circle map is said to be a quasi-symmetry if there exists $K \geqslant 1$ such that for any pair $I, J$ of adjacent and equal intervals, $\operatorname{Leb}(f(I)) \leqslant K \operatorname{Leb}(f(J))$ and vice versa. Suppose that $P$ is a property that subsets of the circle might satisfy. We say that a particular set has quasi-symmetric stability for $P$ if its image under any quasi-symmetry satisfies $P$.

It is of interest to find explicit subsets that display quasi-symmetric stability. By definition, a quasi-symmetry is Hölder continuous; hence it preserves zero Hausdorff dimension. However, there are quasi-symmetries that promote a (Lebesgue) measure zero set to full measure. Specifically, Tukia defines a quasi-symmetry and a set with Hausdorff dimension less than one such that its image by the quasi-symmetry has its complement with Hausdorff dimension less than one. See [21, Main theorem]. By contrast, quasi-symmetries in higher dimensions are quasi-conformal and hence absolutely continuous.

With the circle being the boundary at infinity for the hyperbolic plane, singular measures on the circle (and measure zero sets that correspond to these) arise from various contexts in hyperbolic geometry. A quasi-symmetry of the circle extends to a quasi-isometry of the hyperbolic plane. Conversely, quasi-isometries of the hyperbolic plane give rise naturally to quasi-symmetries of the circle. So it is natural to consider the behaviour of such measures and sets under quasi-symmetries. In particular, finite area complete hyperbolic surfaces/orbifolds with cusps have been used to give explicit measure zero sets that are quasisymmetrically stable for various properties (such as measure zero, porosity, etc.). In Section 1.2, we describe some of these sets and survey known results about them. In this note, we expand the list by considering singular stationary measures arising from random walks on non-uniform lattices in $S L(2, \mathbb{R})=\operatorname{Isom}\left(\mathbb{H}^{2}\right)$ (that is, fundamental groups of finite-area complete hyperbolic surfaces/orbifolds with cusps).

A random walk on a lattice $\Gamma$ in $S L(2, \mathbb{R})=\operatorname{Isom}\left(\mathbb{H}^{2}\right)$ is non-elementary if the semigroup generated by its support contains a pair of independent hyperbolic isometries (with distinct
stable/unstable fixed points on the circle at infinity). Projected in to $\mathbb{H}^{2}$ by the group action, a typical sample path converges to a point on $S^{1}=\partial_{\infty} \mathbb{H}^{2}$. The convergence defines a stationary measure on $S^{1}=\partial_{\infty} \mathbb{H}^{2}$ giving the distribution of sample path limit points.

A lattice $\Gamma$ is non-uniform when the quotient $X=\mathbb{H}^{2} / \Gamma$ is finite area but non-compact and so has cusps. It is well known that stationary measures from a large class of random walks (for example, with finite word metric first moment) on a non-uniform lattice are singular. See [12], [5], [3], [9] for different approaches to understanding the singularity. A lattice is uniform if the quotient is a closed orbifold. A long standing conjecture of Guivarc' h -Kaimanovich-Ledrappier states that stationary measures from finitely supported random walks on a uniform lattice are also singular. See [5]. Here, we restrict to non-uniform lattices.

For a non-uniform lattice $\Gamma$, we fix cusp neighbourhoods in $X=\mathbb{H}^{2} / \Gamma$. The lifts of the cusp neighbourhoods give a horoball packing in $\mathbb{H}^{2}$. We fix a base-point $x$ in $\mathbb{H}^{2}$. Given a horoball $H$ and a hyperbolic ray $\gamma$ from $x$ that intersects $H$, we define the excursion of $\gamma$ in $H$ to be the distance (in the path metric on $\partial H$ ) between the entry and exit points. Let $r \geqslant 0$ be a constant. We define the $r$-shadow of $H$ to be the interval at infinity given by hyperbolic rays from $x$ that have an excursion of size at least $r$ in $H$. We say that a random walk on $\Gamma$ has exponential decay for excursions if for any horoball $H$ and any sample path converging in to the shadow of $H$, the probability that it converges into the $r$-shadow decays exponentially in $r$. The decay holds for finitely supported non-elementary random walks on $\Gamma$. It is expected to be true for a larger class of random walks (for example with finite word metric first moment), but this is as yet unproven. We discuss the more general context for exponential decay in Section 5.4 and sketch a proof for finitely supported random walks.

Our main result below settles a conjecture of McMullen.
Theorem 1.1. Suppose that $\Gamma$ is a non-uniform lattice in $\operatorname{SL}(2, \mathbb{R})$. Suppose that $\mu$ is a nonelementary probability distribution on $\Gamma$. Let $v$ be the stationary measure on $S^{1}$ arising from the $\mu$-random walk. Suppose that $v$ has exponential decay for excursions. Then the pushforward $f^{*} v$ by any quasi-symmetry $f: S^{1} \rightarrow S^{1}$ is singular with respect to the Lebesgue measure on $S^{1}$.

Douady-Earle show that a quasi-symmetry $f: S^{1} \rightarrow S^{1}$ always extends to a bijective quasi-isometry $f$ of $\mathbb{H}^{2}$. See [6, Theorem 2] and also [13, Theorem 3.1] for a strengthening of it. The pushforward measure $f^{*} v$ is then the stationary measure arising from the same random wlak on the group of quasi-isometries resulting from conjugating $\Gamma$ by $f$. Theorem 1.1 is then singularity for stationary measures arising from such groups of quasi-isometries.

Let $S$ be a surface homeomorphic to $\mathbb{H}^{2} / \Gamma$. We may view the quotient $\mathbb{H}^{2} / \Gamma$ as defining a marked conformal structure on $S$ thus giving us a point in the corresponding Teichmüller space $\mathcal{T}(S)$ of $S$. When $f$ is $\Gamma$-equivariant then the conjugate $f^{-1} \Gamma f$ is again a non-uniform lattice in $S L(2, \mathbb{R})$ (isomorphic to $\Gamma$ ). Thus, it can be thought of as a different point of $\mathcal{T}(S)$. By restricting to $\Gamma$-equivariant quasi-symmetries, Theorem 1.1 implies the (already known) singularity of stationary measures given by all of $\mathcal{T}(S)$ from singularity at a single point.

For any finitely supported non-elementary random walk on a uniform lattice $\Gamma$ in $S L(2, \mathbb{R})$, random walk drift in the hyperbolic plane is a proper function on the Teichmüller space of the surface that is homeomorphic to $\mathbb{H}^{2} / \Gamma$. See [2, Theorem A]. It follows that for every such random walk on the abstract group there is a uniform lattice in $S L(2, \mathbb{R})$ such that the stationary measure that arises is singular. If Theorem 1.1 was true for uniform lattices then we would deduce the singularity conjectured by Guivarc'h-Kaimanovich-Ledrappier. For
reasons outlined in Section 1.5, we find it hard to speculate whether Theorem 1.1 holds for uniform lattices and if it does hold whether proving it is as hard as proving singularity itself.
1.2. Quasi-symmetric stability: We list some known examples of quasi-symmetric stability.

- Because quasi-symmetries are Hölder continuous, any set with zero Hausdorff dimension remains so under any quasi-symmetry.
- Any set that is a countable union of uniformly porous sets remains so under any quasi-symmetry. See [22, Theorem 4.2].
- Any set that is strongly winning or absolutely winning remains so under any quasisymmetry. See [17, Theorem 2.2].
1.3. Diophantine sets. Some sets arising from geometric considerations have been analysed extensively for quasi-symmetric stability for properties such as above. The most known example is the set $D$ of Diophantine directions on a finite area complete hyperbolic surface $X$ with cusps. $D$ is the set of directions of hyperbolic rays from a base-point that have bounded excursions in the cusp neighbourhoods of $X$. By ergodicity of the hyperbolic geodesic flow, it follows that $D$ has measure zero. Using the concepts of games, McMullen shows that $D$ is an absolute winning set. McMullen also showed that any absolute winning set remains so under any quasi-symmetry. See [17, Theorem 1.3 and Theorem 2.2]. It is known that winning sets have maximal Hausdorff dimension; so $D$ has Hausdorff dimension one. McMullen also proves that $D$ is countable union of uniformly porous sets. See [17, Proposition 1.5].

Porosity will be the key concept in our proof of Theorem 1.1. So we give a proof sketch of the porosity result for $D$. For a positive integer $m \geqslant 2$, we define $D_{m} \subset D$ consisting of hyperbolic rays whose excursions are all bounded above by $m$. Given a horoball $H$, the $m$ shadow of $H$ is a subinterval of the shadow of $H$ with relative size (in the Lebesgue measure) $\approx 1 / m$. The set $D_{m}$ avoids all these sub-intervals. Using an estimate by Sullivan [20] for the number of horoballs of a certain size in an interval, one consructs a horoball heirarchy where at each level we have that

- the horoball shadows have disjoint interiors and cover a set of full measure; and
- for each horoball of the previous level, the shadows of the current level are contained in the $m$-shadow or in one of the components (in the horoball shadow) of the complement of the $m$-shadow.
This shows that the set $D_{m}$ is $1 / m$-porous. It follows that $D$ is a countable union of porous sets, each uniformly porous.
1.4. Stationary measures from non-uniform lattices. For the stationary measures, a typical ray has unbounded excursions and a countable exhaustion by uniformly porous sets is not possible. Using exponential decay of horoball excursions, we build a horoball heirarchy $\mathcal{H}^{(n)}$ for which there is a constant $c>0$ (that depends only on the random walk) such for a typical ray (for the stationary measure) its excursion in a horoball in $\mathcal{H}^{(n)}$ is eventually (that is, for $n$ large enough depending on the ray) bounded above by $c \log n$. See Proposition 5.7. This implies that eventually a typical ray misses sub-intervals of horoball shadows at all levels of the heirarchy. The relative sizes of the sub-intervals in the corresponding shadows approach zero as the level goes to infinity but quantitatively this happens slowly. This brings us to the notion of (non-uniform) porosity and we show that the support of the stationary measure is contained in a countable union of such porous sets.

By a classical theorem of Ahlfors-Beurling (see [1, Theorem 1]), any quasi-symmetry of the circle is given by integration with respect to a doubling measure. By a theorem of Wu (see
[23, Theorem 1]), non-uniform porous sets with a high enough porosity (slow enough decay of the relative sizes of complementary subintervals/holes) are measure zero for any doubling measure, hence quasi-symmetrically stable. Our excursion upper bounds imply the slow enough decay of the relative sizes of avoided subintervals required in Wu's theorem. Our main theorem, namely Theorem 1.1, is then a direct consequence of Wu's theorem.

Singularity of stationary measures is also known in the following analogous contexts.

- Non-uniform lattices in the isometry group of $\mathbb{H}^{n}$ for $n>2$ (see [18]);
- Mapping class groups of finite type surfaces with negative Euler characteristic (see [7], [10]).
The random walk boundaries in the above contexts are the spheres at infinity $S^{n-1}=\partial_{\infty} \mathbb{H}^{n}$ in the first case, and the Thurston boundaries of the corresponding Teichmüller spaces (see [15]). The dimensions of these boundaries are at least two. Since quasi-symmetries (that is, quasi-conformal maps) in higher dimensions are absolutely continuous, quasi-symmetric stability holds directly. However, the question of porosity is still of interest. We expect that the stationary measure has support contained in a countable union of highly porous (but not uniformly porous) subsets of the boundaries.
1.5. Dimensions of stationary measures: Our countable collection of sets that have high enough porosity is constructed using excursion upper bounds. In particular, (up to excising a countable subset of $S^{1}$ ) their union contains the set $D$ of Diophantine directions. Thus, the Hausdorff dimension of their union is one.

On the other hand, the dimension of a singular stationary measure is expected to be strictly less than one thus giving us a stronger form of singularity. See [5, Conjecture 1.21] and in a more general form in [14]. For uniform lattices, finite first moment for a word metric is the same as finite first moment for the hyperbolic metric. So the dimension gap is conjectured for finitely supported random walks on uniform lattices, the general picture for diffused random walks being unclear.

For non-uniform lattices, finite first moment in the word metric implies finite first moment for the hyperbolic metric but the converse is not true because the parabolic subgroups (corresponding to the cusps) are exponentially distorted. Hence, the dimension gap may be conjectured for more general random walks that have finite word metric first moment. Our countable union of porous sets while serving the purpose for Theorem 1.1 is nevertheless too large to give information regarding the dimension of the stationary measure.

Despite the dimension conjecture for finitely supported random walks on uniform lattices, we think it to be likely that the stationary measure exhibits no porosity. This should rule out our approach in proving Theorem 1.1 for uniform lattices. It is also possible that there exists a non-equivariant quasi-symmetry that does straighten the stationary measure in this case.
1.6. Bounds on maximal excursions of typical geodesics: Similar upper bounds hold when all cusp excursions of a typical hyperbolic ray for the stationary measure are considered. Specifically, there is a constant $c>0$ that depends only on the random walk such that along a typical ray the $n^{\text {th }}$ excursion $E_{n}$ is eventually (that is, for all $n$ large enough depending on the ray) at most $c \log n$. This result follows directly from the exponential decay of excursions and the easy direction of the Borel-Cantelli lemma. In fact, the upper bounds also hold in analogous contexts such as for non-uniform lattices in Isom $\left(\mathbb{H}^{n}\right)$ for $n>2$ and mapping class groups of finite type surfaces with negative Euler characteristic.

Our construction of the horoball hierarchy is more involved because it is built to track porosity. So it is not $\Gamma$-equivariant. Excursions in these horoballs are only a subset of all cusp excursions. Similarity of the bounds suggest that the set of excursions in the horoballs of the heirarchy has positive density among all excursions.

The derivation of a logarithmic lower bound for the maximal excursion over all excursions up to time $T$ is harder and we do not attempt it here. It should require the following two properties.

- A lower bound for the tail distribution (in the stationary measure) of an excursion.
- A quasi-independence of excursions.

We expect these properties to be true when the random walk generates $\Gamma$ as a semi-group and is finitely supported (more generally, has a finite first word metric moment). The arguments are likely to be quite technical. By knowing logarithmic lower bounds we would describe more precisely a full measure set for the stationary measure (in particular, it would not contain $D$ ). It might be sufficient to prove the dimension gap (referred to in Section 1.5).

For the Lebesgue measure, the excursions are much larger. In particular, there is a constant $c^{\prime}>0$ such that for any $c>1$ the largest excursion $E_{\max }(n)$ along any Lebesgue-typical ray is eventually (that is for $n$ large enough depending on the ray) at most $c^{\prime} n(\log n)^{c}$. See [8, Lemma 5.5]. Lower bounds for $E_{\max }(n)$ are also known; for example Sullivan's log-law says that the $\lim \left(\log E_{\max }(n) / \log n\right)$ is $1 / \log 2$. See [20].
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## 2. Horoball heirarchies

Let $\Gamma$ be a non-uniform lattice in $S L(2, \mathbb{R})$. Let $N_{1}, \cdots, N_{k}$ be (pairwise disjoint) cusp neighbourhoods in the finite area hyperbolic surface $X=\mathbb{H} / \Gamma$. Lifting these neighbourhoods to $\mathbb{H}$, let $\mathcal{H}$ be the resulting horoball collection. We call the complement $\mathbb{H}-\cup_{H \in \mathcal{H}} H$ the thick part $\mathcal{T}$. We fix a base-point $x$ in $\mathcal{T}$.

Let $p, p^{\prime}$ be a pair of distinct points in $S^{1}=\partial_{\infty} \mathbb{H}$. They partition $S^{1}$ into a pair of intervals with disjoint interiors. Using the orientation on $S^{1}$ (compatible with orientation on $\mathbb{H}$ ), we define $\left[p, p^{\prime}\right] \subset S^{1}$ to be the interval that is positively oriented as we traverse from $p$ to $p^{\prime}$.

Let $H$ be any horoball in $\mathcal{H}$ with $p_{\infty}(H)$ in $S^{1}$ its point at infinity. Let $\gamma^{-}(H)$ and $\gamma^{+}(H)$ be hyperbolic rays from $x$ with endpoints $p^{-}(H)$ and $p^{+}(H)$ in $S^{1}$ such that

- $\gamma^{-}(H)$ and $\gamma^{+}(H)$ and tangent to $H$; and
- $p_{\infty}(H)$ is contained in $\left[p^{-}(H), p^{+}(H)\right]$.

We call the interval $\mathrm{Sh}_{0}(H)=\left[p^{-}(H), p^{+}(H)\right]$ the shadow at infinity of $H$.
Let $\gamma$ be a geodesic ray from $x$ that converges to a point in $\mathrm{Sh}_{0}(H)$. Let $x_{p}$ and $x_{p}^{\prime}$ be the points of intersection of $\gamma$ with $\partial H$. We define the excursion $E(\gamma, H)$ to be the distance along $\partial H$ between $x_{p}$ and $x_{p}^{\prime}$.

To define shadows more generally, suppose that $r \geqslant 0$. Let $\gamma_{r}^{-}$and $\gamma_{r}^{+}$be geodesic rays from $x$ such that

- $\gamma_{r}^{-}$converges to a point $p_{r}^{-}$in $\left[p^{-}(H), p_{\infty}(H)\right]$;
- $\gamma_{r}^{+}$converges to a point $p_{r}^{+}$in $\left[p_{\infty}(H), p^{+}(H)\right]$; and
- $E\left(\gamma_{r}^{-}, H\right)=r$ and $E\left(\gamma_{r}^{+}, H\right)=r$.

We define the shadow $\mathrm{Sh}_{r}(H)$ to be the interval $\left[p_{r}^{-}, p_{r}^{+}\right.$]. By definition, if $r<s$ then $\mathrm{Sh}_{s}(H)$ is nested in $\mathrm{Sh}_{r}(H)$ that is, the closure of $\mathrm{Sh}_{s}(H)$ is contained in the interior of $\mathrm{Sh}_{r}(H)$. Let Leb denote the Lebesgue measure on $S^{1}$. It follows from the nesting (and basic hyperbolic geometry) that there is a constant $c>0$ such that for any horoball $H$ in $\mathcal{H}$

$$
\begin{equation*}
\operatorname{Leb}\left(\operatorname{Sh}_{r}(H)\right) \geqslant \frac{c}{r} \operatorname{Leb}\left(\operatorname{Sh}_{0}(H)\right) \tag{2.1}
\end{equation*}
$$

See [8, Section 3].
We organise horoballs by the (Euclidean) sizes of their shadows as follows. Suppose that $0<\rho<1$ is a positive constant. We say that a horoball has $\rho$ size $n$ if $\rho^{n+1}<\operatorname{Leb}\left(\operatorname{Sh}_{0}(H)\right) \leqslant$ $\rho^{n}$. Suppose that $I$ is an interval in $S^{1}$. We define $\mathcal{H}_{n}(I)$ to be the collection of those horoballs $H$ in $\mathcal{H}$ such that

- $\mathrm{Sh}_{0}(H)$ is contained in $I$; and
- the $\rho$-size of $H$ is $n$.

We recall [20, Proposition 4]. The proof uses the mixing of the hyperbolic geodesic flow on $T^{1} \mathbb{H}^{2} / \Gamma$. The proposition is not stated as precisely in [20] as it is stated here. However, our statement is implicit in the proof of [20, Proposition 4] requiring no modification other than tracking explicitly how the counting depends on the compact subset of the boundary.
Proposition 2.2. Suppose that $\Gamma$ is a non-uniform lattice in $S L(2, \mathbb{R})$. There are positive constants $C$ and $\rho<1$ such that for any interval I in $S^{1}$ there is a positive integer $n_{I}$ such that

$$
\left|\mathcal{H}_{n}(I)\right| \geqslant C \frac{\operatorname{Leb}(I)}{\rho^{n}}
$$

for all $n \geqslant n_{I}$.
Suppose that $I$ is an interval in $S^{1}$. We define the subcollection $\mathcal{H}_{n}^{\prime}(I) \subset \mathcal{H}_{n}(I)$ by requiring that for any pair of horoballs $H_{1}$ and $H_{2}^{\prime}$ in $\mathcal{H}_{n}^{\prime}(I)$ the shadows $\mathrm{Sh}_{0}\left(H_{1}\right)$ and $\mathrm{Sh}_{0}\left(H_{2}\right)$ have disjoint interiors. We use Proposition 2.2 to deduce:
Lemma 2.3. Suppose that $\Gamma$ is a non-uniform lattice in $\operatorname{SL}(2, \mathbb{R})$. There is a positive constant $0<$ $\delta<1$ such that for any interval I in $S^{1}$ there is a positive integer $n$ such that the union $U=$ $\cup_{H \in \mathcal{H}_{n}^{\prime}(I)} \mathrm{Sh}_{0}(H)$ satisfies

$$
\operatorname{Leb}(U)>\delta \operatorname{Leb}(I)
$$

Proof. By Proposition 2.2, we may fix $n_{1}=n_{I}$ and consider horoballs $H$ such that $\mathrm{Sh}_{0}(H)$ is contained in $I$ and the $\rho$-size of $H$ is $n_{1}$. By definition, each horoball shadow can intersect at most $1 / \rho$ shadows of size $n$. Thus, we may retain at least $C^{\prime} \operatorname{Leb}(I) / \rho^{n}$ horoballs (where $\left.C^{\prime}=c \rho\right)$ to get the collection $\mathcal{H}_{n}^{\prime}(I)$. It follows that for $U=\cup_{H \in \mathcal{H}_{n}^{\prime}(I)} \mathrm{Sh}_{0}(H)$ we have

$$
\operatorname{Leb}(U)=\sum_{H \in \mathcal{H}_{n}^{\prime}(I)} \operatorname{Sh}_{0}(H) \geqslant\left(\left|\mathcal{H}_{n}^{\prime}(I)\right|\right) \cdot \rho^{n+1}=C \rho^{2} \operatorname{Leb}(I)
$$

The lemma follows by choosing $\delta<C \rho^{2}$.
We use Lemma 2.3 to deduce:
Lemma 2.4. Suppose that $I$ is an interval in $S^{1}$. There exists a collection $\mathcal{H}(I) \subset \mathcal{H}$ of horoballs $H$ such that

- for distinct horoballs $H$ and $H^{\prime}$ in $\mathcal{H}(I)$, the interiors of $\mathrm{Sh}_{0}(H)$ and $\mathrm{Sh}_{0}\left(H^{\prime}\right)$ are disjoint; and

Proof. By Lemma 2.3 applied to $I$, we can find a collection $\mathcal{H}(I, 1)=\mathcal{H}_{n}^{\prime}(I)$ of horoballs all of $\rho$-size $n$ such that their shadows have disjoint interiors and their union $U(I, 1)$ satisfies $\operatorname{Leb}(U(I, 1))>\delta \operatorname{Leb}(I)$. We now consider the intervals that are the components of $I-$ $U(I, 1)$. These are finitely many and again by Lemma 2.3 applied to each of these intervals, we can find a collection of horoballs for each interval such that their shadows have disjoint interiors and their union has proportion (for the Lebesgue measure) at least $\delta$ in the interval. We may then define $\mathcal{H}(I, 2)$ to be the collection of these horoballs and define $U(I, 2)$ to be the union of their shadows. Note that $\operatorname{Leb}(U(I, 2))>\delta \operatorname{Leb}(I-U(I, 1))$.

Iterating infinitely, we get a sequence horoball collections $\mathcal{H}(I, k)$ and subsets $U(I, k)$ of $I$ such that

- $U(I, 0)$ is the empty set,
- $U(I, k) \subset I-\cup_{i \leqslant k-1} U(I, i)$;
- for $k \geqslant 1$ we have that $U(I, k)$ is the finite union of shadows of horoballs in $\mathcal{H}(I, k)$ with pairwise disjoint interiors, and
- there is a constant $\delta>0$ such that $\operatorname{Leb}(U(I, k))>\delta \operatorname{Leb}\left(I-\cup_{i \leqslant k-1} U(I, i)\right)$ for all $k \geqslant 1$.
It follows that the infinite union $\cup_{k} U(I, k)$ is a countable union of horoball shadows with pairwise disjoint interiors and it has full measure in I.

Using Lemma 2.4, we define the first sub-collection $\mathcal{H}^{(1)}$ of horoballs in our hierarchy.
Lemma 2.5. There exists a collection $\mathcal{H}^{(1)} \subset \mathcal{H}$ of horoballs $H$ such that

- for distinct horoballs $H$ and $H^{\prime}$ in $\mathcal{H}^{(1)}$, the interiors of $\mathrm{Sh}_{0}(H)$ and $\mathrm{Sh}_{0}\left(H^{\prime}\right)$ are disjoint; and

Proof. The proof follows directly from Lemma 2.4 applied to $S^{1}$.

Let $\alpha>0$ be a constant. The horoball sub-collections of $\mathcal{H}$ that we define below depend on $\alpha$. We drop the dependence from our notation to keep it simple.

We define the sequence $r_{n}=\alpha \log n$. We now define the collection $\mathcal{H}^{(n)}$ for $n \geqslant 2$ inductively. Let $H$ be a horoball in $\mathcal{H}^{(n-1)}$. Let $\mathcal{K}$ be the collection of horoballs $H^{\prime}$ such that $\mathrm{Sh}_{0}\left(H^{\prime}\right)$ is contained in either $\mathrm{Sh}_{r_{n}}(H)$ or in the closure of one of the components of $\mathrm{Sh}_{0}(H)-\mathrm{Sh}_{r_{n}}(H)$. By Lemma 2.4, we can find a collection $\mathcal{H}_{H}^{(n)}$ of horoballs in $\mathcal{K}$ such that

- for distinct horoballs $K$ and $K^{\prime}$ in $\mathcal{H}_{H}^{(n)}$, the interiors of $\mathrm{Sh}_{0}(K)$ and $\mathrm{Sh}_{0}\left(K^{\prime}\right)$ are disjoint; and
- the union $\bigcup_{K \in \mathcal{H}_{H}^{(n)}} \mathrm{Sh}_{0}(K)$ has full measure in $\mathrm{Sh}_{0}(H)$.

See Figure 2.6.
We define $\mathcal{L}_{H}^{(n)}$ to be the collection of those horoballs $H^{\prime}$ in $\mathcal{H}_{H}^{(n)}$ such that $\mathrm{Sh}_{0}\left(H^{\prime}\right)$ is contained in $\mathrm{Sh}_{r_{n}}(H)$. We define $L_{H}^{(n)}$ to be the union $\bigcup_{H^{\prime} \in \mathcal{L}_{H}^{(n)}} \mathrm{Sh}_{0}\left(H^{\prime}\right)$. By Equation (2.1),

$$
\operatorname{Leb}\left(L_{H}^{(n)}\right) \geqslant \frac{c}{r_{n}} \operatorname{Leb}\left(\operatorname{Sh}_{0}(H)\right)
$$



Figure 2.6. Illustration of how the collection $\mathcal{H}_{H}^{(n)}$ is constructed.
We define

$$
\mathcal{H}^{(n)}=\bigcup_{H \in \mathcal{H}^{(n-1)}} \mathcal{H}_{H}^{(n)}
$$

Note that the union $U^{(n)}=\bigcup_{H \in \mathcal{H}^{(n)}} \operatorname{Sh}(H)$ is full Leb-measure in $S^{1}$.
We also define

$$
\mathcal{L}^{(n)}=\bigcup_{H \in \mathcal{H}^{(n-1)}} \mathcal{L}_{H}^{(n)}
$$

We define $L^{(n)}$ to be the union $\bigcup_{H \in \mathcal{L}^{(n)}} \mathrm{Sh}_{0}(H)$. Since $U^{(n)}$ has full measure, by Equation (2.1) we get

$$
\operatorname{Leb}\left(L^{(n)}\right) \geqslant \frac{c}{r_{n}} \operatorname{Leb}\left(U^{(n)}\right)=\frac{c}{r_{n}}
$$

## 3. Porous Sets

We fix $n \geqslant 1$. We define the set $Y(n, n+1)$ by

$$
\begin{equation*}
Y(n, n+1)=U^{(n)}-L^{(n)} \tag{3.1}
\end{equation*}
$$

Let $m>n$. We define the set $Y_{n, m}$ inductively by $Y(n, m)=Y(n, m-1)-L^{(m-1)} \cap Y(n, m-$ 1). We define $X_{n}$ to be the infinite intersection $\bigcap_{m>n} Y(n, m)$.

Definition 3.2. Let $a_{n}$ be a sequence of numbers satisfying $0<a_{n}<1$. A set $X$ is said to be $\left\{a_{n}\right\}$ porous if there is a sequence of coverings $\mathcal{E}^{(n)}=\{E(n, j)\}$ by intervals with disjoint interiors such that

- $E(n, j)-X$ contains an interval $E^{\prime}(n, j)$ such that $\operatorname{Leb}\left(E^{\prime}(n, j) \geqslant a_{n} \operatorname{Leb}(E(n, j))\right.$,
- $\cup_{\mathcal{E}^{(n+1)}} E(n+1, k)$ is contained in $\bigcup_{\mathcal{E}^{(n)}} E(n, j)-E^{\prime}(n, j)$, and
- $\sup _{\mathcal{E}^{(n)}} \operatorname{Leb}(E(n, j))$ tends to 0 as $n \rightarrow \infty$.

We call the intervals $E^{\prime}(n, j)$ the holes at level $j$.

Proposition 3.3. The set $X_{n}$ is porous for the sequence

$$
a_{j}=\frac{c}{r_{n+j}}
$$

Proof. For each $j \geqslant 1$, we set $\mathcal{E}^{(j)}$ to be collection of shadows $\mathrm{Sh}_{0}(H)$ where $H$ is contained in $Y(n, n+j-1)$. For each such shadow, we set the hole to be $\mathrm{Sh}_{r_{n+j}}(H)$. By Equation (2.1), the required estimate for the relative measure follows.

## 4. QUASI-SYMMETRIES AND MEASURE ZERO SETS

We say that two intervals $I$ and $J$ in $S^{1}$ are adjacent and equal if they share an endpoint and have equal length.

Definition 4.1. A map $f: S^{1} \rightarrow S^{1}$ is said to be a quasi-symmetry if there exists a constant $K \geqslant 1$ such that for any adjacent and equal intervals I and J in $S^{1}$ we have

$$
\frac{1}{K} \leqslant \frac{\operatorname{Leb}(f(I))}{\operatorname{Leb}(f(J))} \leqslant K
$$

Definition 4.2. A measure $\mu$ on $S^{1}$ is said to be a doubling measure if there exists a constant $M>0$ such that for any $p \in S^{1}$ and any radius $R>0$, the interval $I(p, 2 R)$ centred at $p$ and with radius $2 R$ satisfies

$$
0<\mu(I(p, 2 R))<M \mu(I(p, R))
$$

By a classical result of Ahlfors-Beurling, namely [1, Theorem 1], a map $f: S^{1} \rightarrow S^{1}$ is quasi-symmetric if and only if it is given by integration with respect to a doubling measure, that is, if there exists a doubling measure $\mu$ and $\theta>0$ such that

$$
f\left(p^{\prime}\right)-f(p)=\int_{p}^{p^{\prime}} d \mu
$$

We begin by recalling [23, Theorem 1]
Theorem 4.3. Suppose that $a_{n}$ is a sequence of numbers satisfying $0<a_{n}<1$. Suppose that for all real numbers $\kappa \geqslant 1$ the sum

$$
\sum_{n=1}^{\infty}\left(a_{n}\right)^{\kappa}
$$

diverges. Then any $\left\{a_{n}\right\}$-porous set has measure zero for any doubling measure.
In other words, if the sequence $\left\{a_{n}\right\}$ satisfies the hypothesis then for any quasi-symmetric map $f: S^{1} \rightarrow S^{1}$ the image of an $\left\{a_{n}\right\}$-porous set has measure zero for the Lebesgue measure.

Recall our definition of the set $X_{n}$. By Theorem 4.3, we deduce:
Proposition 4.4. Suppose that $f: S^{1} \rightarrow S^{1}$ is a quasi-symmetry. We have

$$
\operatorname{Leb}\left(f\left(X_{n}\right)\right)=0
$$

Proof. Note that for any $\kappa \geqslant 1$, the sequence $a_{j}$ in Proposition 3.3 satisfies

$$
\sum_{j} a_{j}^{\kappa}=\sum_{j}\left(\frac{c}{\alpha \log (n+j)}\right)^{\kappa}=\infty
$$

The proposition follows from Theorem 4.3.

## 5. Stationary measures and horoball excursions

5.1. Horoball excursions and nested sets in the thick part. The action of the (virtually free) group $\Gamma$ on the thick part $\mathcal{T}$ is co-compact. By the Švarc-Milnor lemma, the thick part $\mathcal{T}$ with the induced path metric $D$ is quasi-isometric to $\Gamma$ for any word metric on $\Gamma$ defined by a finite symmetric generating set. Let $\pi: \mathbb{H} \rightarrow \mathcal{T}$ be the closest point projection. We call the projection $\pi(\gamma)$ of a hyperbolic geodesic $\gamma$ a projected path.

Let $x$ be a point in $\mathcal{T}$ and $H$ a horoball with the point at infinity $p_{\infty}(H) \in S^{1}$. With $x$ as a base-point, let $\gamma_{\infty}(H)$ be the geodesic ray from $x$ to $p_{\infty}(H)$. Recall that $\gamma^{-}(H)$ and $\gamma^{+}(H)$ are rays from $x$ tangent to $H$. Also recall that $\gamma_{r}^{-}$and $\gamma_{r}^{+}$are rays from $x$ that have excursions of size $r$ in $H$.

Let $s_{r}^{-}$denote the arc $\pi\left(\gamma_{r}^{-}\right) \cap \partial H$. Let $x_{1}^{-}$and $x_{2}^{-}$be the endpoints of $s_{r}^{-}$with $x_{r}^{-}$the farthest (among the two) from $x$. For any geodesic ray $\gamma$ from $x$ such that $\gamma_{r}^{-}, \gamma, \gamma_{\infty}$ are in the counter-clockwise order, the projected path $\pi(\gamma)$ contains the arc $s_{r}^{-}$. For such a ray, let $t_{r}^{-}$ denote the initial segment of $\pi(\gamma)$ from $x$ to $x_{2}^{-}$. Let $R^{-}(r)$ be the set of such rays and define the thick shadow

$$
C^{-}(x, H ; r)=\bigcup_{\gamma \in R^{-}(r)} \pi(\gamma)-t_{r}^{-}
$$

We similarly define $s_{r}^{+}$and use it to define the thick shadow $C^{+}(x, H ; r)$ using the ray $\gamma_{r}^{+}$.
We define $\partial_{\infty} C^{-}(x, H ; r)$ to be the set of endpoints at infinity of the rays in $R^{-}(r)$. It follows that $\partial_{\infty} C^{-}(x, H ; r)=\left[p_{r}^{-}(H), p_{\infty}(H)\right]$ and $\partial_{\infty}=\left[p_{\infty}(H), p_{r}^{+}(H)\right]$ and so their union is $\mathrm{Sh}_{r}(H)$.

Lemma 5.2. There exists constants $K \geqslant 1$ and $r_{0}>0$ that depend only on $\Gamma$ and the choice of cusp neighbourhoods such that for all $r$ large enough and $\left(r^{\prime}-r\right) / K-r_{0}>0$, we have

$$
\begin{equation*}
D\left(\mathcal{T}-C^{-}(x, H ; r), C^{-}\left(x, H ; r^{\prime}\right)\right)=\frac{r^{\prime}-r}{K}-r_{0}, \tag{5.3}
\end{equation*}
$$

Proof. By [9, Lemma 2.1], a projected path is a quasi-geodesic in $\mathcal{T}$. Since $C^{-}(x, H ; R)$ is defined with a projected arc as a prefix, the lemma follows.
5.4. Exponential decay for cusp excursions. Let $\mu$ be a non-elementary probability distribution on $\Gamma$. We consider the random walk on $\Gamma$ given by $\mu$. We call the random product $\omega_{n}=g_{1} \cdots g_{n}$ (where each $g_{i}$ is sampled independently by $\mu$ ) a sample path of length $n$. Sample paths of length $n$ are distributed by the $n$-fold convolution $\mu^{(n)}$. The path space for the random walk is the space $\Gamma^{\mathbb{N}}$ of one-sided infinite sequence in $\Gamma$ equipped with the pushforward of the product measure under the map $\left(g_{1}, g_{2}, \cdots\right) \rightarrow\left(\omega_{1}, \omega_{2}, \cdots\right)$. We denote this measure as $v$ and it has marginals $\mu^{(n)}$. The group $\Gamma$ acts on the left while the random walk increments $g_{n}$ act by right multiplication $\omega_{n-1} \rightarrow \omega_{n}$.

Let $x$ be any base-point in $\mathbb{H}$. For $v$-almost every infinite sample path $\omega^{+}$, the sequence $\omega_{n}^{+} x$ in $H$ converges to a point $p=p(\omega)$ in $S^{1}$. We may then pushforward $v$ to get a measure on $S^{1}$ that gives the distribution of $p$. In a slight abuse of notation, we continue to denote the measure on $S^{1}$ by $v$.

We now define what it means for a random walk to have exponential decay for excursions.

Definition 5.5. A random walk has exponential decay for excursions if there exists constants $c \geqslant 0$ and $\alpha>0$ such that for any horoball $H$ in $\mathcal{H}$ and all constants $r, s \geqslant 0$ such that $s \geqslant r+c$, we have

$$
v\left(\operatorname{Sh}_{s}(H)\right) \leqslant e^{-\alpha} v\left(\operatorname{Sh}_{r}(H)\right)
$$

Note that the exponential decay implies that for $r \geqslant 2 c$ and $\beta=\alpha / 2 c$

$$
v\left(\operatorname{Sh}_{r}(H)\right) \leqslant e^{-\beta r} v\left(\operatorname{Sh}_{0}(H)\right)
$$

For completeness, we include a proof sketch that finitely supported random walks that generate the lattice as a semi-group have the desired exponential decay for excursions. As a special case of more general exponential decay of shadows (see [16, Lemma 2.10]), the exponential decay is true for any finitely supported non-elementary random walk. In the specific contexts mentioned here, we expect the exponential decay to be true for all nonelementary random walks with finite word metric first moment, but this is open at present.

Proof sketch for finitely supported random walks that generate $\Gamma$ : For any pair of points $z \neq z^{\prime}$ in $\mathcal{T}$, let $H\left(z, z^{\prime}\right)$ be the half-space (in the metric $D$ on $\mathcal{T}$ ) of points closer to $z^{\prime}$ than $z$. Let $\partial_{\infty} H\left(z, z^{\prime}\right)$ be the limit set at infinity of $H\left(z, z^{\prime}\right)$.

Suppose that $H$ is a horoball. Suppose that $\xi>0$ is a constant. Suppose that $r^{\prime}=r+$ $3 \xi$. Recall that $x \in \mathcal{T}$ is our base-point. We consider the thick shadows $C^{-}\left(x, H ; r^{\prime}\right) \subset$ $C^{-}(x, H ; r)$ and similarly $C^{+}\left(x, H ; r^{\prime}\right) \subset C^{+}(x, H ; r)$. We define the set $C^{-}(x, H ; r+\xi)-$ $C^{-}(x, H ; r+2 \xi)$ to be the midpoint set for the pair $C^{-}\left(x, H ; r^{\prime}\right) \subset C^{-}(x, H ; r)$. We denote the set by $M^{-}(\xi)$. Similarly, we define the set $C^{+}(x, H ; r+\xi)-C^{+}(x, H ; r+2 \xi)$ to be the midpoint set for the pair $C^{+}\left(x, H ; r^{\prime}\right) \subset C^{+}(x, H ; r)$. We denote it by $M^{+}(\xi)$.

By Lemma 5.2 , there is $\xi^{\prime}>0$ such that for any $\xi^{\prime} \geqslant \xi^{\prime}$ there is $u=u(\xi)>0$ such that such that for any point $x^{\prime}$ in the midpoint set $M^{-}(\xi)$ (similarly $\left.M^{+}(\xi)\right)$ there is a point $x^{\prime \prime}$ with $D\left(x^{\prime}, x^{\prime \prime}\right)=u$ such that $\partial_{\infty} H\left(x^{\prime}, x^{\prime \prime}\right)$ is contained in the interior of $\left[p_{r}^{-}(H), p_{r^{\prime}}^{-}(H)\right]$ (similarly $\left[p_{r^{\prime}}^{+}(H), p_{r}^{+}(H)\right]$.

Since the random walk generates $\Gamma$, there is a constant $\kappa>0$ such that for any $z$ such that $D(x, z)=u$, we have $v\left(\partial_{\infty} H(x, z)\right)>\kappa$.

Let $\delta$ be maximum of $d_{\mathcal{T}}\left(x_{0}, g x_{0}\right)$ over $g$ in support of $\mu$. Suppose that $H$ is a horoball. Recall the constant $r_{0}$ from Equation (5.3). Let $\xi=\max \left\{\xi^{\prime}, K\left(\delta+r_{0}\right)\right\}$. Suppose that $r^{\prime}>r$ be such that $r^{\prime}=r+3 \xi$. Then any sample path that converges to $\mathrm{Sh}_{r^{\prime}}(H)$ intersects $M^{-}(\xi)$ or $M^{+}(\xi)$. Conditioning on the point of intersection, we conclude that the probability for a sample path starting at this point to converge in to $\mathrm{Sh}_{r}(H)-\mathrm{Sh}_{r^{\prime}}(H)$ is at least $\kappa$. The claimed decay for $v$ now follows.
5.6. Upper bounds for excursions in the horoball heirarchy. We now return to the horoball hierarchies $\mathcal{H}^{(n)}$ constructed in Section 2. Suppose that $p$ lies in $S^{1}$. Suppose that $\gamma$ is the hyperbolic ray from the base point $x$ to $p$. If $\gamma$ has an excursion in a horoball $H$ in $\mathcal{H}^{(n)}$ then we denote the excursion by $E^{(n)}(\gamma)$. By definition, for $n \geqslant 2$, the set $L^{(n)}$ is the set of points $p$ in $S^{1}$ such that the ray $\gamma$ from $x$ to $p$ satisfies $E^{(n-1)}(\gamma) \geqslant r_{n}$.

We prove:

Proposition 5.7. There exist a constant $\alpha>0$ that depends only on $\mu$ such that for the horoball collections $\mathcal{H}^{(n)}$ and for $v$-almost every $p \in S^{1}$ there is a positive integer $N=N(\gamma)$ such that with $r_{n}=\alpha \log n$, we have

$$
E^{(n-1)}(\gamma)<r_{n}
$$

for all $n>N$.
Proof. We fix $\alpha>0$ to satisfy $\alpha>1 / \beta$. We then consider horoball heirarchies $\mathcal{H}^{(n)}$ corresponding to the sequence $r_{n}=\alpha \log n$. By the exponential decay of excursion from Definition 5.5 , it follows that the

$$
v\left(L^{(n)}\right)<e^{-\beta r_{n}}=e^{-\beta \alpha \log n}=\frac{1}{n^{\beta \alpha}}
$$

It follows that

$$
\sum_{n \geqslant 2} v\left(L^{(n)}\right)<\infty
$$

By the Borel-Cantelli lemma (the easy direction), it follows that $v\left(\lim \sup L^{(n)}\right)=0$. The conclusion follows.

We deduce:
Corollary 5.8. We have

$$
v\left(\bigcup_{n \geqslant 1} X_{n}\right)=1
$$

Proof. Suppose that $p$ is $v$-typical and suppose that $\gamma$ is the hyperbolic ray from $x$ to $p$. By Proposition 5.7, there exists a positive number $N$ such that $E^{(n-1)}(\gamma)<r_{n}$ for all $n>N$. Thus, $p$ lies in $X_{N}$ and we get our conclusion.

Proof of Theorem 1.1. Set $X=\cup_{n \geqslant 1} X_{n}$. Suppose that $f: S^{1} \rightarrow S^{1}$ is a quasi-symmetry. Consider $f(X)$. By Proposition 4.4, we have

$$
\operatorname{Leb}(f(X)) \leqslant \sum_{n \geqslant 1} \operatorname{Leb}\left(f\left(X_{n}\right)\right)=0
$$

On the other hand, by definition of a pushforward measure and Corollary 5.8

$$
f^{*} v(f(X))=v(X)=1
$$

Thus $f^{*} v$ is singular, as desired.

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School of Mathematics and Statistics, University of Glasgow, University Place, Glasgow, G12 8QQ United Kingdom<br>Email address: Aitor.Azemar@glasgow.ac.uk

School of Mathematics and Statistics, University of Glasgow, University Place, Glasgow, G12
8QQ United Kingdom
Email address: Vaibhav.Gadre@glasgow.ac.uk

