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# On the effects of $NH_3$ addition to a reacting mixture of $H_2/CH_4$ under MILD combustion regime: Numerical modeling with a modified EDC combustion model

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### ABSTRACT

This paper examines the behavior of reacting NH<sub>3</sub>/H<sub>2</sub>/CH<sub>4</sub> mixtures in moderate or intense low oxygen dilution (MILD) condition. A series of axisymmetric, turbulent reacting flow simulations are carried out incorporating a modified version of eddy dissipation concept and a few reaction mechanisms. The effects of adding a progressively increasing amount of NH3 to a reacting H2/CH4 mixture in moderate condition are investigated. It is observed that addition of NH<sub>3</sub> to MILD combustion leads to markedly different behaviors compared to that in conventional combustion. Most notably, the inherently strong preheating of reactants in MILD combustion causes thermal cracking of NH<sub>3</sub> prior to ignition. The resultant production of H<sub>2</sub> profoundly affects the reacting flow as such increasing the NH<sub>3</sub> mass fraction in the fuel blend decreases the flame lift-off. Further, unlike that in conventional combustion, adding NH<sub>3</sub> to MILD combustion increases the process reactivity. In addition, the usual flame thickening typically seen in NH<sub>3</sub> flames is not observed here, while in keeping with the thermodynamic predictions, NH<sub>3</sub> addition lowers the temperature of combustion products. The results also show that in sharp contrast to that reported for conventional combustion, addition of NH<sub>3</sub> in MILD condition does not increase the emission of NO, while the mass fraction of NO2 drops slightly. Overall, it is concluded that MILD combustion could be a promising route to NH<sub>3</sub> combustion.

# 1. Introduction

Ammonia (NH<sub>3</sub>) as a carbon-free fuel has recently attracted considerable interest as the fuel of future with significant economic advantages [1]. This flammable inorganic substance is deemed to play a major role in decarbonization of the energy sector and development of H<sub>2</sub> economy. NH3 consists of nearly 18% H2 by mass and is recognized as contributing to H<sub>2</sub> combustion through a well-developed and economically viable transportation and storage infrastructure [1]. In contrast to pure H<sub>2</sub>, NH<sub>3</sub> is transported relatively easily in liquid form. Also, the high octane number (about 130) of NH<sub>3</sub> helps eliminate knocking in internal combustion engines [2]. However, currently burning of NH<sub>3</sub> faces important challenges, which should be addressed before its extensive use in combustion systems. Amongst these, is the low

combustion efficiency of NH<sub>3</sub> burning systems [2,3], originating from the low reactivity of NH3 compared to hydrocarbons. Further, flame instabilities and high emissions of nitrogen oxides (NO<sub>x</sub>) are observed during combustion of  $NH_3$  [4-6]. The instabilities are chiefly due to the low flame speed of NH<sub>3</sub> [7,8], which leads to the flame blow-off and extinction [9]. Formation of NO in NH<sub>3</sub> combustion is primarily along the pathway of fuel NO and is much higher under lean conditions, requiring the employment of expensive post-treatment systems to reduce  $NO_x$  emissions [10,11]. However, while  $NH_3$  can be burned in fuel-rich conditions to minimize emissions, this combustion regime demands removal of unburned NH<sub>3</sub> and other pollutants (N<sub>2</sub>O) from the exhaust gases [10,11]. The issues of low flame velocity of pure NH<sub>3</sub> (~6-8 cm/s under stoichiometric conditions [12]) and high ignition energy [4] can be addressed by mixing NH<sub>3</sub> with other fuels to increase the reactivity of the unburned mixture [13]. Combustion of NH<sub>3</sub> and H<sub>2</sub>

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Nomenclature		$ au^*$	fluid mean residence time
		$ au_C$	chemical time scale
$C_{\tau}$	fine structure volume constant	$ au_\eta$	Kolmogorov mixing time scale
$C_{\gamma}$	residence time constant	ε	turbulent dissipation rate
$CD_1, CD_2$	2 coefficients of ECM	ν	Kinematic viscosity
$Da_{\eta}$	Damköhler number at the Kolmogorov scale	$\dot{\omega}_i$	mean reaction rate of species i
L*	L* length scale		ations
Re <sub>T</sub>	turbulent Reynolds number	EDC	addy dissinction concent
u*	turbulent reacting fine structure characteristic velocity	EDC	eddy dissipation concept
$S_T$	turbulent flame speed	MILD	moderate or intense low oxygen dilution
$S_L$	laminar flame speed	ECM	energy cascade model
Y <sub>i</sub>	mass fraction of the i <sup>th</sup>	Supersci	ripts
Greek let	tors	bar	time-averaged values
	thermal diffusivity	tilde	Favre-averaged values
u ~	Turbulant thormal diffusivity	*	Characteristics quantities
$\alpha_{\rm T}$			1
$\gamma_{\lambda}$	fine structures mass fraction		

is an attractive way of achieving zero-carbon combustion while improving combustion performance (flammability and stability).

In a spark-ignition engine, Mørchet al. [14] examined the combustion of NH<sub>3</sub>/H<sub>2</sub> mixtures. They discovered that a high H<sub>2</sub> mass fraction in the mixture and an excess air ratio range between 1.1 and 1.4 resulted in greater NO<sub>x</sub> emissions. It was also found that NO<sub>x</sub> levels could be reduced to those of gasoline combustion through post-combustion destruction, implying that NH<sub>3</sub>/H<sub>2</sub> mixtures could be an alternative to hydrocarbons in spark ignition engine. An analysis of burning NH<sub>3</sub>/H<sub>2</sub> blended fuels in a porous burner was carried out by Nozari et al. [15]. The porous burners were found capable of burning fuel blends with a significant fraction of NH<sub>3</sub>. In a swirl burner, Valera-Medina et al. [16] studies combustion of NH<sub>3</sub>/H<sub>2</sub> and NH<sub>3</sub>/CH<sub>4</sub> mixtures. These authors discovered that the NH3/H2 fuel blends had a limited equivalence ratio ranges for which the flames stayed stable, and that NO<sub>x</sub> emission might overtake the maximum levels of 4500 ppm. The flame stability was attained for a wider equivalence ratio ranges in the NH<sub>3</sub>/CH<sub>4</sub> fuel mixtures, whereas NO<sub>x</sub> emission might reach a maximum levels of 2600 ppm near stoichiometry. In a gas turbine, Kurata et al. [11] investigated combustion of NH3 as well as NH3/CH4 blends. It was reported that the NOx emission for pure NH3 combustion was the lowest, barely around 1000 ppm. Moreover, they hypothesized that the separation of rich and lean zones inside the burner caused an interaction between unburned NH<sub>3</sub> and NO<sub>x</sub>, lowering the overall NO<sub>x</sub> emissions. Consequently, however, a large NH<sub>3</sub> concentration (more than1500 ppm) emerged at the combustion chambers outlet. The burning of NH<sub>3</sub>/CH<sub>4</sub> fuel mixtures, on the other hand, resulted in significantly lower NH<sub>3</sub> emissions. Rocha et al. [17] described a study of NH<sub>3</sub>/H<sub>2</sub>/air ignitions, premixed flame propagations, and NO emissions using chemical kinetic modeling. They showed that NH<sub>3</sub> combustion could feature long ignition delays and low flame speeds and that adding H<sub>2</sub> to the process could enhance the flame speed significantly with the expense of major increases in NO<sub>x</sub> emissions.

Han et al. [18] studied the laminar firing rates of NH<sub>3</sub>/air, NH<sub>3</sub>/H<sub>2</sub>/ air, NH<sub>3</sub>/CO/air, and NH<sub>3</sub>/CH<sub>4</sub>/air combustions by the heat flux methodology to better understand the properties of NH<sub>3</sub> combustion. Blending NH<sub>3</sub> with H<sub>2</sub> has been proven to be an effective way of boosting the flame speed of NH<sub>3</sub>-based fuel mixtures. Ramos et al. [19] investigated NH<sub>3</sub>, CO, and NO emissions from premixed NH<sub>3</sub>/CH<sub>4</sub>/air flames. They discovered that NO<sub>x</sub> emissions increased with the NH<sub>3</sub> molar fraction up to 0.5 before they diminish. Furthermore, NO<sub>x</sub> emissions dropped when the equivalence ratio decreased toward fuel-lean settings, but CO and NH<sub>3</sub> emission remained relatively low regardless of the operating conditions. Analyses of premixed swirl flames fed with NH<sub>3</sub>/ H<sub>2</sub> blends under very-lean to stoichiometric conditions was reported by Zhu et al. [20]. These authors showed that the OH\* chemiluminescence level could be utilized as a proxy for the NO mole fraction under certain conditions. Based on the weak flame responses recorded in a micro flow reactor with a controlled temperature profiles, Murakami et al. [21] evaluate the influence of blending and equivalence ratio on the oxidation and reactivity of dimethyl ether/NH<sub>3</sub> blends. As the NH<sub>3</sub> mass fraction in fuel mixture increased from 0 to 15%, dimethyl ether oxidation was boosted and mixture reactivity increased under stoichiometric condition. CO oxidation was inhibited when the NH<sub>3</sub> mass fraction increased from 15 to 50%, and the mixed reactivity were significantly reduced. At intermediate temperatures of 800–1000 K, NO<sub>x</sub> generated by NH<sub>3</sub> oxidation enhanced the reactivity of radicals via the NO–NO<sub>2</sub> catalytic loop.

The other practical approach to using  $NH_3$  as a fuel for power generation is the application of the well-demonstrated concept of MILD combustion. This offers the unique potential to increase combustion performance while at the same time reducing pollutant emissions [22,23]. In general, MILD regime includes preheating and dilution, which boost flame stability and decrease thermal  $NO_x$  emission [22]. The concept of MILD combustion shares some of its fundamental features with modern engine technologies, such as homogeneous charge compression ignition and exhaust gas recirculation. It should be noted that MILD combustion technology can adapt to the complexity and variability of power consumption patterns while avoiding drastic changes in existing power plants [24].

Some recent studies have shown the potential for using carbon-free fuels such as NH<sub>3</sub> or NH<sub>3</sub>/H<sub>2</sub> mixtures with promising results for MILD combustion. For instance, Manna et al. [25] experimentally characterized NH<sub>3</sub> oxidation and pyrolysis processes in model reactors. In the jet stirred flow reactor, the data obtained allowed for the identification of three distinct kinetic regimes in NH<sub>3</sub> oxidation: low temperatures (T<sub>in</sub> < 1130 K), middle temperatures (1130 < T<sub>in</sub> < 1250 K), and high temperatures (T<sub>in</sub> greater than 1250 K). Surface heterogeneous reactions appeared to have no influence on NH<sub>3</sub> reactivity and H<sub>2</sub> profiles, but they did have an effect on NO<sub>x</sub> concentration in the low-intermediate temperature range. Sue et al. [26] experimentally showed that NH<sub>3</sub>-MILD combustion decreases NO<sub>x</sub> emissions by a substantial amount. Sabia et al. [27] reported that NH<sub>3</sub>/CH<sub>4</sub> mixtures are important when biogas is used as a fuel, and MILD combustion could represent an important solution to keep NO<sub>x</sub> emissions very low.

Ferrarotti et al. [28] developed an innovative, multidisciplinary technique for analyzing  $NH_3/H_2$  mixtures in flameless combustion. They demonstrated that adding a trace quantity of  $NH_3$  to pure  $H_2$  significantly boosted NO emissions (10% by volume). When operating at sufficiently close to stoichiometric circumstances (equivalence ratio of

0.95), an optimal trade-off between NO<sub>x</sub> emission and NH<sub>3</sub> slip was discovered. According to these authors, enlarging the air injector could cut down the emissions. When the equivalence ratio of 0.8 was reached, emissions were reduced further more. Sorrentino et al. [29] studied the stability and emission features of the NH3-MILD regime for several equivalence ratios and inlet temperatures. They reported lower NO<sub>x</sub> levels close to the stoichiometric condition and stable combustion at above 1300 K. Ariemma et al. [30] examined the behavior of H<sub>2</sub>Oassisted NH3-MILD combustion under both premixed and non-premixed conditions. Their findings showed that H<sub>2</sub>O addition to the combusting flow in the NH3-MILD regime is a very simple and efficient way of reducing NO emissions in the fuel-lean MILD regime. Rocha et al. [31] conducted numerical simulations of an idealized combustor and observed that MILD combustion could produce minimal NO<sub>x</sub> (50 ppm) and unburned NH<sub>3</sub> (0.1 ppm) emissions. The authors did highlight, however, achieving this requires a large amount of exhaust gas recirculation, which might be unfeasible. Sorrentino et al. [32] analyzed the reactive structure of a one-dimensional steady NH3-counter flow diffusion flame. These authors indicated that in comparison with the case of fuel preheating, NH<sub>3</sub>-MILD combustion is more stable at lower fuel dilution levels when the oxidant flow is preheated. Ariemma et al. [33] concentrated on the NH<sub>3</sub>/CH<sub>4</sub>-MILD combustion in a lab-scale burner in order to examine the gaseous pollutant emissions and stability of the process limitations as a function of the equivalence ratio and the NH<sub>3</sub>/ CH<sub>4</sub> ratio. It was shown that compared to pure NH<sub>3</sub>, the use of NH<sub>3</sub>/CH<sub>4</sub> mixtures could make the system much more stable in terms of working temperatures and equivalence ratios. However, utilization of fuel blends resulted in greater NO<sub>x</sub> emissions as compared to those emitted by pure NH<sub>3</sub> and CH<sub>4</sub>. Zhao et al. [34] explored the role of co-flow O<sub>2</sub> and temperature on MILD combustion and fuel-NO mechanism of CH<sub>4</sub>/NH<sub>3</sub> jet diffusion flames in hot O2/CO2 co-flow. Temperature and co-flow O2 were varied from 3% to 30% and 1300 K to 2100 K, respectively. The results indicated that when co-flow O2 was 21%, MILD-oxy combustion could be fully achieved for all co-flow temperatures. However, a higher co-flow temperature was required to maintain MILD regime when coflow O2 was 24%. Sun et al. [35] examined the homogeneous oxidation of CH<sub>4</sub>/NH<sub>3</sub>/NO/NO<sub>2</sub> mixtures at temperatures ranging from 600 K to 1200 K in a laminar flow reactor under atmospheric pressure. The results indicated that when NH<sub>3</sub>/NO<sub>x</sub> mixtures were oxidized, the NO/  $NO_2$  ratio affected the conversion of NO via the reaction  $NH_2 + NO_2 =$ H<sub>2</sub>NO + NO. The inclusion of CH<sub>4</sub> considerably increased the oxidation reactivity of the  $NH_3/NO_x$  combination. As the amount of  $CH_4$  added rose, the generation of H<sub>2</sub>O increased while the synthesis of N<sub>2</sub> dropped. Furthermore, when CH<sub>4</sub>/NH<sub>3</sub>/NO<sub>x</sub> mixtures were oxidized, the NO<sub>2</sub>/NO ratio had little influence on the generation of main products and intermediates. Manna et al. [36] studied the oxidation of highly diluted NH<sub>3</sub>-H<sub>2</sub> mixtures at low-intermediate temperatures in a jet stirred flow reactor. In a parametric study, they varied the mass fraction of H<sub>2</sub> in the fuel blend, the mixture equivalence ratio, and the inlet temperatures. It was observed that H<sub>2</sub> addition barely improved the system reactivity at low-intermediate temperatures and pushed lower reactivity and dynamic regimes toward lower reactor temperatures.

Studies on NH<sub>3</sub> combustion in conventional conditions have already revealed the requirement for highly swirling burners to increase stability and minimize nitrogen oxide emissions. Therefore, MILD combustion might be considered as the asymptotic solution to the problem. Literature indicates that gains in performance in terms of efficiency and flexibility could emerge. Nevertheless, a rigorous study on this topic has not been reported yet. Therefore, the focus of the current work is on NH<sub>3</sub>/CH<sub>4</sub>/H<sub>2</sub> mixtures, with relatively high mole fraction of CH<sub>4</sub>, to gain a profound understanding of the effects of NH<sub>3</sub> on the main parameters of MILD combustion. These include preheating zone length, flame front, the thickness of the reacting zone as well as NO<sub>x</sub> and CO emissions. A Dally's burner, like the one used in earlier studies [37,38], was used for the current investigation. In general, the difficulty of simulating this regime has been one of the key challenges, stemming from the uncommon interactions between transport and chemistry in MILD combustion [39,40]. To resolve this issue, several modeling approaches have already been proposed [41-45]. According to the flow conditions in the present work and findings of the previous studies, the method presented by Cuoci et al. [46,47] seems to be the most appropriate. Hence, in the proceeding study the extended edcSimpleSMOKE solver, proposed by this group is employed with different reaction mechanisms and k- $\varepsilon$  turbulence model.

Numerical method.

# 1.1. Eddy dissipation concept

The eddy dissipation concept (EDC) [23] is useful for the prediction of turbulent reacting flows, especially in the modes in which combustion kinetics play an important role (e.g. MILD combustion) [48,49]. An energy cascade model (ECM) provides the fine structures mass fraction ( $\gamma_{\lambda}$ ) and the fluid mean residence time ( $\tau^*$ ) [32]. To calculate the residence time, there are two constants including the fine structure volume and residence time constants. These are set to 2.137 and 0.4083 by default. Based on the classical ECM, the fine structure scale is of the same order of the Kolmogorov scale as follows;

$$Re^* = \frac{u^*L^*}{\nu} = \frac{2}{3} \frac{C_{D2}}{C_{D1}} = 2.5 \tag{1}$$

It should be noted that this model is appropriate for high Reynolds number flows in which there is a significant distinction amongst turbulent scales. There is no apparent distinction between turbulence on large and small scales in MILD conditions, and reactions occur across a wide range of scales [47]. Consequently, the chemical reaction proceeds in a thick reaction region that corresponds to the integral length scale. This modifies the characteristic scale of the reaction structure because energy is transferred at a higher frequency than that of the reaction structure in the spectrum [50]. Hence, it is essential to consider the nature of MILD combustion mode and modify the cascade model. This is also needed to explain how the ECM parameters are affected by the flow and response qualities represented by the Reynolds and Damköhler numbers. Dilution and preheating of oxidizer under MILD condition create a unique distributed reaction area [47]. Damköhler number about one characterizes the system evolving towards the perfectly mixed condition. Parente et al. [47] modified the classical cascade model by assuming that MILD condition occurs in the so-called 'Distributed Reaction Regime' with small-scale and high-intensity turbulence. They noted that the use of Eq. (15) could be a good first-order approximation for the characteristic velocities of the reaction structure for MILD condition. Further, it could be utilized to derive the dependency of the coefficients of ECM  $(CD_1 \text{ and } CD_2)$  on the dimensionless number of combusting flows.

$$S_T \approx S_L \sqrt{\frac{\alpha_T + \alpha}{\alpha}} \approx S_L \sqrt{Re_T + 1}$$
 (2)

The relation between the turbulent dissipation rate and turbulent reacting fine structure characteristic velocity  $(u^*)$  indicates.

$$\varepsilon \propto C_{D2} \nu u^{*2} / L^{*2} \tag{3}$$

where  $u^* \approx S_T \approx S_L \sqrt{Re_T + 1}$ . Combination of Eqs. (2) and (3) leads to the following definition;

$$\varepsilon \propto C_{D2} \nu \frac{u^{*2}}{L^{*2}} = C_{D2} \nu \frac{S_L^2 (Re_T + 1)}{L^{*2}}$$
(4)

in which,  $L^*/S_L$  is the characteristic chemical time scale ( $\tau_C$ ) of the reaction structure according to the classical treatment of a turbulent premixed flame [36]. Therefore,  $\tau_C$  is expressed as a function of the Kolmogorov mixing time scale ( $\tau_\eta$ ) by the Damköhler number of reacting flow. Thus, the dissipation is defined as:

$$\varepsilon \alpha \frac{[C_{D2}\nu(Re_T+1)]}{\tau_c^2} = \frac{\left[C_{D2}\nu(Re_T+1)Da_{\eta}^2\right]}{\tau_{\eta}^2}$$
(5)

Here,  $Da_{\eta} = \tau_{\eta}/\tau_c$  is Damköhler number which has a value of one in MILD condition. Based on $\tau_{\eta} = (\nu/\epsilon)^{\frac{1}{2}}$ , the dependency of C<sub>D2</sub> on *Re<sub>T</sub>* and  $Da_{\eta}$  is obtained as below

$$C_{D2} \propto \frac{1}{\left[Da_{\eta}^{2}(Re_{T}+1)\right]}$$
(6)

Also,

$$C_{\tau} \propto \frac{1}{Da_{\eta} \sqrt{Re_{T} + 1}} \tag{7}$$

A similar method for determination of the dependency of  $C_{\gamma}$  on  $Re_T$  and  $Da_{\eta}$  leads to the following definition

$$C_{\gamma} = \left(\frac{3C_{D2}}{4C_{D1}^2}\right)^{\frac{1}{4}} \propto \left[\frac{(Re_T+1)}{C_{D2}}\right]^{\frac{1}{4}} \propto Da_{\eta}^{\frac{1}{2}} (Re_T+1)^{\frac{1}{2}}$$
(8)

# 1.2. Plug flow reactor

Disregarding mixing term between the fine structures and their surroundings is a straightforward technique to eliminate the instabilities and reducing the computational burden [51,52]. Since back-mixing is neglected in the species conservation equations and only chemistry is integrated over time, this technique is referred to as the plug flow reactor approach. A plug flow reactor is a one-dimensional, steady-state reactor that has both inflow and outflow, as well as a distribution of characteristics along the flow direction. In this technique, reactions occur across time scales dictated by Arrhenius rates. When the flow velocity, area, and pressure are constant, the governing equations may be translated into transient equations, such as the species mass balance, which can be written as:

$$\frac{dY_k}{dt} = \frac{R_k^*}{\rho^*} \tag{9}$$

The initial conditions are considered as the current species and temperature in the cell meaning that Eq. (9) is integrated from  $Y_k(t = 0) = \tilde{Y}_k$  to  $Y_k(t = \tau^*) = Y_k^*$ .

Numerical settings, configuration, computational grid and boundary conditions.

The OpenSMOKE library [46] and the extended edcSimpleSMOKE solver [53] with a simple algorithm in a two-dimensional axisymmetric configuration were used in this study to detect the physics of MILD combustion in several working conditions. Table 1 shows the details of the present numerical set-up. In the present work, the turbulent Prandtl number value was set to one. The existing experimental data [54] for a typical Jet-in-Hot-Coflow burner (Dally burner) with the configuration shown in Fig. 1 was used to validate the results of the extended EDC model [54,55]. Dally's burner contains a hot co-flow and a fuel inlet nozzle with the diameters of 80 and 4.25 mm, respectively. It should be

# Table 1

Numerical setup for the simulations in the present work.

Code	OpenFOAM-7
Solver	Modified edcSimpleSMOKE
Combustion model	EDC
Multicomponent diffusion	yes
Molecular viscosity	Sutherland law
Coupling of Pressure and velocity fields	SIMPLE
Schemes of discretization	second order
Fine structure reactor	Plug fine structure
Turbulent Schmidt number	0.7
Turbulent Prandtl number	1.0
mechanism of chemical reaction	GRI-Mech 3.0

noted that the burner is equipped with a wind tunnel with a diameter of 420 mm. In the tunnel, air flows at a speed of 3.2 m/s in a direction parallel to the burner axis [56]. Table 2 indicates the working conditions considered in the current simulations. For the current MILD combustion problem, H<sub>2</sub>, CH<sub>4</sub>, N<sub>2</sub>, H<sub>2</sub>O, and CO<sub>2</sub> are the main components of the inlet flows. Furthermore, H<sub>2</sub>/CH<sub>4</sub> fuel is injected with the inlet Reynolds number of 10,000. It is noted that the velocity inlet and pressure outlet boundary conditions were set for the upstream and downstream, respectively.

It has been already shown that the turbulence intensity of the hot coflow and inlet nozzles did not affect the solution [37,38,55,57,58]. Yet, the fuel inlet turbulence intensity is important [38]. Christo et al. [55] observed that the mean turbulent kinetic energy of  $60 \text{ m}^2/\text{s}^2$  at the fuel inlet yields the best agreement between the simulations and experimental data for the distribution of species mass fraction and temperature.

Furthermore, the grid sensitivity analysis is depicted in Fig. 2. Three cell configurations of 11000, 33000, and 99,000 cells elements were employed to assess the grid independency of the numerical solutions. The results for the mass fractions of  $CH_2O$  and OH radicals are shown in Fig. 2. According to the data in this figure, the relative differences between the results of 33,000 and 99,000 elements are approximately 0.2%, indicating that the results are independent of the grid for the cell configurations with more than 33,000 elements. As a result, this grid was employed for the rest of the investigations reported in this study.

Moreover, as shown in Fig. 1, a finer mesh around the fuel and coflow inlets was employed, with high grid resolution in the reaction region to moderate the computational cost. Note that, in Fig. 1, there are 6 lines by the axial distances of 5 (L1), 10 (L2), 15 (L3), 20 (L4), 25 (L5), and 30 (L6) cm to compare different cases.

## 1.3. Validation

This section presents validation of the numerical simulations by comparing the detected radial mass fraction with the experimentally determined temperature distributions [54] at the axial distances (ADs) of 30 and 120 mm. For this purpose, two different approaches, including the EDC method with default values and modified-EDC with modification of  $C_{D1}$  and  $C_{D2}$  coefficients as in Eqs. (1)–(8), were used. Fig. 3 illustrates a comparison between the predicted results and the experimental data [54] for the radial distributions of OH, CO, and H<sub>2</sub>O as well as those of temperature at the ADs of 30 and 120 mm from the inlet. The default-EDC model is clearly incapable of forecasting the distribution of chemical species and the maximum temperature, as this model could not detect the local extinctions and predict the characteristics of MILD combustion like reaction zone thickness. However, due to the physical properties of MILD combustion, this model has a large error in predicting CO and OH species, as also reported by other researchers [59,60]. This is due to the cooling and extinction influences of the premixed reacting flow inside the secondary combustion chamber where the behavior of the reacting flow deviates from the MILD condition [55,60]. In general, due to the importance of these two species in combustion process, the use of an EDC model is not recommended. However, by changing CD1 and CD2 coefficients to 0.167 and 1, satisfactory agreement was found between the simulations and the experimental results. The increase in C<sub>D1</sub> and C<sub>D2</sub> coefficients enhances the time coefficient of fine structure. This means that, considering the wider reaction zones, smoother gradients exist to reduce the driving forces and lower the temperature with the higher dilution. The characteristics of MILD combustion were recently analyzed by Minamoto et al. [34], who found that the reaction region in this regime was distributed over a considerable part of the burner. Further, the interactions amongst the reactive regions both distribute these zones and uniform the temperature distribution. In this case, u' has a significant effect on the combustion process, and Kolmogorov scales have the ability to enter and widen the preheating region, as well as the reacting area, resulting in a



Fig. 1. Schematic view of computational domain and cell configuration.

Table 2	
Working conditions considered in the simulations	[54]

Inlet condition	Fuel	Co-flow	wind Tunnel
Composition, mass fraction Temperature [K]	$\begin{array}{l} H_2 = 11.1, \\ CH_4 = 89.9 \\ 300 \end{array}$	$\begin{array}{l} O_2 = 9,  N_2 = 79,  H_2 O \\ = 6.5,  CO_2 = 5.5 \\ 1300 \end{array}$	$\begin{array}{l} O_2 = 23.3, N_2 \\ = 76.7 \\ 300 \end{array}$

thicker and diffused flame. [61]. By increasing the role of u',  $Re_T$  increases, leading to an increase in the fine structure coefficient (with a lower percentage than the fine structure time coefficient). Considering these parameters and characteristics, the modified-EDC model can accurately predict the structure of MILD combustion in both centerline of the burner and radial distance in various ADs. In addition, the following equation [62] was used for calculating the error of the present work in predicting the experimental data. Using Eq. (10) shows that the average values of numerical results accuracy in predicting H<sub>2</sub>O, CO, OH, and temperature are 96.5, 97.3, 97.3, and 97.6%, respectively.

$$R^{2} = \frac{\sum_{i=1}^{n} \left( Z_{\exp,i} - Z_{simulation,mean} \right)^{2} - \sum_{i=1}^{n} \left( Z_{\exp,i} - Z_{simulation,i} \right)^{2}}{\sum_{i=1}^{n} \left( Z_{\exp,i} - Z_{simulation,mean} \right)^{2}}$$
(10)

Next, the capability of the GRI-Mech 3.0 in predicting the behavior of MILD combustion with  $NH_3$  added to the fuel blend is investigated. There exist several reaction mechanisms which include  $NH_3$  [63,64]. Importantly, however, most of them are restricted to the mixtures of  $H_2$  and  $NH_3$ . Further, the chemical mechanisms of  $CH_4/H_2/NH_3$  by Okafor et al. [65],  $NH_3/H_2/CH_4$ -R1 [60], and  $NH_3/H_2/CH_4$ -R2 [60], by the characteristics shown in Table 3, have been used. The fuel compositions in this section include  $H_2 = 11.1$ ,  $CH_4 = 69.9$ , and  $NH_3 = 20$  (NH20 case in Table 4), with the inlet velocity and temperature, as well as the inlet co-flow and wind tunnel characteristics as those described in Table 2. Fig. 4 shows a comparison among the outcomes of these four mechanisms for different species and the process temperature at different axial locations. It is observed that using the GRI-Mech 3.0 mechanism, similar results are obtained to those of other mechanisms. This is because the mass fraction of CH<sub>4</sub> in the investigated blend is much greater than that



Fig. 2. The grid independency tests.



Fig. 3. Comparison between the current simulations and the experimental data of Ref. [54].

# Table 3 The main characteristics of the reaction mechanisms used in the present work.

Reaction mechanism	Okafor et al. [65] (M1)	GRI-Mech 3.0 [66] (M2)	NH <sub>3</sub> /H <sub>2</sub> / CH <sub>4</sub> -R2 [60] (M3)	NH <sub>3</sub> /H <sub>2</sub> / CH <sub>4</sub> -R1 [60] (M4)
Number of reactions	130	325	420	634
Number of species	42	53	51	74

Mass fraction of the inlet fuel composition $(H_2/G_2)$	CH₄/N	IH3)	١.

Inlet condition	Mass fractions of $\mathrm{H}_2, \mathrm{CH}_4$ and $\mathrm{NH}_3$ in the reactants
NH0	$H_2 = 11.1, CH_4 = 88.9, NH_3 = 0$
NH5	$H_2 = 11.1, CH_4 = 83.9, NH_3 = 5$
NH10	${\rm H}_2=11.1,{\rm CH}_4=79.9,{\rm NH}_3=10$
NH15	$H_2 = 11.1, CH_4 = 73.9, NH_3 = 15$
NH20	$\rm H_2 = 11.1,  CH_4 = 69.9,  NH_3 = 20$

of NH<sub>3</sub>. It was observed that when the molar fraction of NH<sub>3</sub> was pushed greater than 30% and that of CH<sub>4</sub> was lower than 50%, the results of GRI-Mech 3.0 were quite different from the results of the other NH<sub>3</sub> mechanisms. Furthermore, it is noted that the majority of reactions in the Okafor mechanism [64] are from the GRI-Mech 3.0 mechanism. It follows that the use of the GRI-Mech 3.0 mechanism is acceptable for the blends with low concentration of NH<sub>3</sub>.

## 2. Results and discussions

This section presents the effects of adding NH<sub>3</sub> with various mass fractions on the behavior of MILD combustion in the configuration

shown in Fig. 1. To this end, four different cases shown in Table 4 were considered. It is noted that the other inlet conditions such as inlet temperatures, velocities, and mass fractions, are the same as those presented in Table 2.

Fig. 5 shows the preheating zone in which  $T \leq T_{\text{self ignition}},$  before reaching the maximum H<sub>2</sub> and CO mass fraction [61], and no reactions take place, for MILD combustion of NH<sub>3</sub>/CH<sub>4</sub>/H<sub>2</sub>. This figure is drawn in the direction of the starting point of the reaction, which is higher than the centerline, and it includes the results for different mass fractions of NH<sub>3</sub> in the fuel blend. Due to the low reactivity and burning velocity of NH<sub>3</sub>, thicker flames are often formed in conventional combustion. Per unit mass, NH<sub>3</sub> has a much lower enthalpy of combustion compared to methane. Therefore, compared to CH<sub>4</sub>, NH<sub>3</sub> combustion is projected to generate lower temperatures. However, as Fig. 6 shows, H<sub>2</sub> is produced by the thermal decomposition of NH<sub>3</sub> in the zone near the inlet nozzle. Evidently, increasing the mass fraction of NH<sub>3</sub> in the fuel component leads to an increase in the H2 mass fraction before the axial distance of  $0.07\ m.$  Beyond this point, the burning rate increases because of the enhancement of H<sub>2</sub>, and the mass fraction of H<sub>2</sub> drops by increasing the mass fraction of NH<sub>3</sub>. This is because the preheating temperature in MILD combustion exceeds that needed for the thermal cracking of NH<sub>3</sub> (NH<sub>3</sub> thermal cracking occurs at 1200 K [67,68]). As a result, the ignition delay is reduced while the decomposition of NH3 increases the reaction rate. In addition, blending NH3 with H2 and CH4 results in an enhancement of flame speed and process heat release. Therefore, using NH<sub>3</sub> in MILD condition shortens the preheating zone, as shown in Fig. 5. Based on this figure, as the mass fraction of NH<sub>3</sub> increases from 0 to 20%, the preheating length reduces by up to 85%. However, it should be noted that the preheating zones for NH15 and NH20 are approximately the same length. This behavior is further illustrated in Fig. 7, in which the distribution of OH mass fraction has been plotted as a function of NH<sub>3</sub> mass fraction. The OH free radical is an indicator of heat release from the combustion process and it also marks the reacting zone [69]. Evidently,



Fig. 4. Comparison among the results of Okafor et al. (M1), GRI-Mech 3.0 (M2), NH<sub>3</sub>/H<sub>2</sub>/CH<sub>4</sub>-R2 (M3), and NH<sub>3</sub>/H<sub>2</sub>/CH<sub>4</sub>-R1 (M4) mechanisms.

increasing the  $NH_3$  mass fraction in the blended fuels moves the starting point of OH radical production toward to the inlet nozzle, and unlike that in conventional combustion, the addition of  $NH_3$  to MILD combustion increases the process reactivity.

At various axial distances, Fig. 8 depicts the radial variation of OH mass fraction as a flame marker. The low laminar flame speed of  $NH_3$  contributes to the increases in the thickness of the reaction zone. However, based on this figure increasing the  $NH_3$  mass fraction in the blended fuel decreases the thickness of the reacting zone in MILD condition. There are two reasons for such behavior, 1) preheating temperature higher than the thermal cracking temperature of  $NH_3$ , and 2) blending  $NH_3$  with  $CH_4$  and  $H_2$ , which increases the flame speed with respect to the  $NH_3$  combustion without  $H_2/CH_4$ . When comparing the amount of OH at various dilution levels, there is a competition between the rate at which the temperature falls when  $NH_3$  is added and the rate at which H is added when ammonia is cracked. In this regard, Fig. 8 shows that the localized OH mass fraction is reduced by up to 8.3% in the NH20

case in comparison with NH0. This is attributed to the reduction of temperature in NH20, which is about 80 K compared with NH0 (Fig. 11) due to an overall increase in the nitrogen content of the mixture. Further, adding  $NH_3$  and reduction of  $CH_4$  in the fuel component increase the overall H production in the mixture due to cracking (Fig. 15) which is the other reason for the reduction of OH radical. This confirms that the flame becomes weaker with increasing the  $NH_3$  mass fraction.

In general, the phenomena related to decreases in the OH free radical have been explained as a "weakening" of the flame front [70]. The weakening of the reaction zone is primarily attributted to the rate of heat release, which is marginally impacted by the high temperature of the coflow and shear layer mixing. Fig. 9 demonstrates the zone of maximum mass fraction of OH for various values for the cases stated in Table 4. The yellow color in the cadres shows the maximum value of the OH radical. Evidently, increasing the mass fraction of  $CH_4$ , results in a reduction in the area and thickness of these zones. Since OH radical



Fig. 5. Preheating zone length as a function of the NH<sub>3</sub> mass fraction (x-axis is normalized with the diameter of the inlet fuel nozzle).



Fig. 6. Radial distribution of  $H_2$  mass fraction in the axial distances of a- 0.01, b- 0.03, c-0.05, d = 0.07, e = 0.09, and f = 0.11 m (x-axis is normalized with the diameter of the inlet fuel nozzle).

acts as a flame marker, the weakened reaction area increases as the mass fraction of  $\rm NH_3$  increases in the zones far from the inlet.

In addition, changes in the fuel components affect the flow properties at the downstream of the reacting flow. One of these effects is related to changes in flow temperature, which lead to variation in the local density and viscosity of the reacting flow. These two properties change the kinematic viscosity of reacting flow. Local variations in flow density also change the local velocity through the mass conservation law, resulting in



Fig. 7. Distribution of OH radical (starting point of OH production) for the cases shown in Table 4.

changes in the local flow dissipation rate. Followed by changes in local dissipation rate and kinematic viscosity, the local residence time of reacting flow is changed. Fig. 10 shows variations in the local ratio of kinematic viscosity and local turbulent dissipation rate as an index of the local flow residence time. It is evident that an increase in the mass fraction of  $NH_3$  followed by a reduction of  $CH_4$  in the fuel component leads to an enhancement of local residence time. It should be noted that combustion temperature depends on heat release, which has nonlinear behavior. Therefore, the changes in the local residence time should be nonlinear.

Fig. 11 represents the radial temperature and CO<sub>2</sub> mole fraction distribution for different mass fractions of NH<sub>3</sub>. As shown in this figure, increasing the mass fraction of NH<sub>3</sub> along with decreasing the volume fraction of CH4 decreases the process temperature. The maximum temperature of CH<sub>4</sub>/H<sub>2</sub> combustion is about 60-80 K greater than that of NH<sub>3</sub>/CH<sub>4</sub>/H<sub>2</sub> (NH20) combustions under the same inlet condition (inlet velocity and temperature). This is due to enthalpy for the combustion of NH3 being significantly less than that of CH4, which results in less heat release per unit mass of the fuel burned. An important reason for the drop in the process temperature is the increase in nitrogen mass fraction due to addion of NH<sub>3</sub>. Furthermore, changes in fuel components leads to changes in the heat capacity of the mixture as an important factor to change in process temperature. In addition, the drop in CO<sub>2</sub> mass fraction by reducing the methane content of the inlet fuel mixture decreases the overall emissivity of the reacting flow in the combustion process leading to a reduction of the radiation heat transfer. These explain the differences in the maximum process temperature at various axial distances. Nonetheless, the differences among the maximum temperatures at all axial distances of the computational domain were relatively small in MILD condition.

So far, it has been demonstrated that  $NH_3/H_2/CH_4$  mixtures have lower flame temperatures (Fig. 11) and heat release rates than  $H_2/CH_4$  mixtures (Fig. 8), and that the disparities increase as the NH<sub>3</sub> mass fraction increases. Due to the reduction of CH<sub>4</sub> mass fraction in the mixture brought about by an increase in NH<sub>3</sub> mass fraction, the carbon concentration falls, resulting in more available oxygen per carbon and a reduction of CO emissions by up to 15.3%. (see Fig. 12). In addition, the addition of NH<sub>3</sub> increases the mixture's residence time, allowing for more time for the mixture to react. Increasing the residence time by increasing the NH<sub>3</sub> mass fraction boosts the H<sub>2</sub>O mass fraction production as a measure of combustion efficiency at a given axial location (Fig. 11) while adding NH<sub>3</sub> followed by reduction of CH<sub>4</sub> reduces H atoms, H<sub>2</sub>O increases, meaning that the process is more completed. Therefore, adding NH<sub>3</sub> improves combustion efficiency and completeness.

In NH<sub>3</sub> flames, NO is produced mainly through the fuel NO pathway, and a higher NO concentration is anticipated in the reaction zone [71]. Fig. 13 depicts the radial distribution of NO and NO<sub>2</sub> mass fractions. According to Fig. 13, increasing the NH<sub>3</sub> mass fraction reduces the amount of NO2 at the downstream at the axial distance of 0.2 m. This figure shows that by addition of NH<sub>3</sub>, the amount of NO rapidly increases near the inlet nozzles. For instance, by adding only 5% of NH<sub>3</sub> mass fraction to the CH<sub>4</sub>/H<sub>2</sub> blend, NO emission raised by several orders of magnitude reaching the unacceptable levels of 1000 ppm in the region near to the inlet nozzles. In general, the HNO intermediate channel is the major NO generation pathway in combustion of NH<sub>3</sub> [72]. Fig. 14 shows the dependency of HNO mass fraction to enhancement of NH<sub>3</sub> in the fuel blend. Addition of NH<sub>3</sub> enhances the HNO mass fraction sharply at the zone near the inlet nozzles with a trend similar to NO changes. However, in the upstream and after the axial distance of 0.3 m, both NO and HNO emissions become the same for all investigated cases. The process temperature could be reason for the existence of same trend in NO and HNO production for various cases downstream. According to Fig. 11, the downstream temperature is nearly equal for all cases,



Fig. 8. Radial distribution of OH mass fraction in axial locations of L1 to L6 (x- axis is normalized with inlet fuel nozzle diameter).



Fig. 9. The zone of maximum OH mass fraction in various cases showing the weakened reaction area (yellow color in cadre shows the OH maximum zone). (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

resulting in similar NO production in all cases. In addition, the analyses show that because of dilution, the smaller  $O_2$  mass fraction in the NH<sub>3</sub>-MILD combustion leads to lower radical pool concentrations. Fig. 15 illustrates the radial distributions of the H/O mass fraction at different axial distances. Addition of NH<sub>3</sub> to the fuel blend reduces the concentration of both radicals slightly. It is known that H abstraction consumes NH<sub>3</sub> primary through the reaction with OH [1]. Other secondary consumption processes involve reactions with H and O, the most typical result of which is NH<sub>2</sub>. Dependent mostly on concentration of the O/H radicals, NH<sub>i</sub> (i = 0–2) oxidation can produce NO (through an HNO

intermediate) or reduce NO (by NH<sub>i</sub> + NO reactions) [73]. Therefore, the slight reduction of O/H radicals prevents NH<sub>i</sub> radicals from increasing NO levels at the downstream. Since O/H radicals are reduced in the NH<sub>3</sub>/CH<sub>4</sub>/H<sub>2</sub> blends, and an increase in NH<sub>3</sub> further decreases these radicals in MILD regime, mixing of NH<sub>3</sub> with CH<sub>4</sub> and H<sub>2</sub> has no significant effect on NO concentration. In addition, the reduction of poor radicals inhibit the path NH<sub>2</sub>  $\rightarrow$  HNO  $\rightarrow$  NO through NH<sub>2</sub> + O<=>HNO + H which is the primary source of HNO [1]. Simultaneously, the reaction of NH<sub>2</sub> + NO<=>N<sub>2</sub> + H<sub>2</sub>O and the path of NH<sub>2</sub>  $\rightarrow$  NNH  $\rightarrow$  N<sub>2</sub> are intensified because the small O<sub>2</sub> mass fraction prevents the increase of



**Fig. 10.** Distribution of the ratio of local kinematic viscosity to local turbulent dissipation rate (x-axis is normalized with the diameter of the inlet fuel nozzle).

NO [26]. Furthermore, the analysis illustrates that the most important reaction for NO<sub>2</sub> generation is NO + HO<sub>2</sub> = NO<sub>2</sub> + OH. Since HO<sub>2</sub> is a common pre-ignition species that normally emerges in the low temperature region, this could explain the reduced NO<sub>2</sub> emission in high-temperature zones (see Fig. 13). It follows that MILD combustion technology could be a promising way of utilizing NH<sub>3</sub> as a carbon free fuel.

# 3. Conclusions

Numerical simulation of MILD combustion of  $NH_3/H_2/CH_4$  blends, with relatively low concentration of ammonia, was carried out using k- $\epsilon$ turbulence model and modified EDC. Adequacy of the employed chemical mechanism and validity of the simulations were demonstrated. The effects of adding varying amounts of  $NH_3$  to MILD combustion were then explored. The key findings of this study can be summarized as follows:

- When NH<sub>3</sub> is added to the MILD combustion regime, it significantly shortens the preheating zone and increases the reacting flow residence time, resulting in more complete combustion.
- The maximum temperature of CH<sub>4</sub>/H<sub>2</sub> (NH0) combustion is approximately 60–80 K higher than that of NH<sub>3</sub>/CH<sub>4</sub>/H<sub>2</sub> (NH20).
- The addition of NH<sub>3</sub> decreases the CO emissions considerably and increases the NO at the inlet area. However, in the downstream, these emissions are roughly similar to those of H<sub>2</sub>-CH<sub>4</sub> MILD combustion.
- Increasing the NH<sub>3</sub> mass fraction in the blended fuels decreases the flame lift-off and, unlike in conventional combustion, the addition of NH<sub>3</sub> to MILD combustion enhances the reactivity.
- Addition of NH<sub>3</sub> results in shrinkage of the zone of maximum mass fraction of OH, so the reaction zone is weaker.
- Adding NH<sub>3</sub> to CH<sub>4</sub>/H<sub>2</sub> blend reduces the thickness of the reaction zone.

Finally, this work showed that MILD combustion could offer a promising route towards achieving clean combustion of ammonia. Further studies, particularly experimental investigations, are necessary to fully clarify this point.



Fig. 11. Radial temperature and  $CO_2$  mole fraction distributions for different cases in axial locations of L1 to L6 (x- axis is normalized with inlet fuel nozzle diameter).



Fig. 12. Radial distributions of CO and H<sub>2</sub>O mass fraction in axial locations of L1 to L6 (x- axis is normalized with inlet fuel nozzle diameter).



Fig. 13. The radial distribution of NO and NO2 mass fractions in axial locations of L1 to L6 (x- axis is normalized with inlet fuel nozzle diameter).

# CRediT authorship contribution statement

**Seyed Mahmood Mousavi:** Conceptualization, Methodology, Validation, Formal analysis, Investigation, Writing – original draft. **Freshteh Sotoudeh:** Conceptualization, Methodology, Validation, Formal

analysis, Investigation, Writing – original draft. **Daeyoung Jun:** Software, Validation. **Bok Jik Lee:** Conceptualization, Methodology, Resources, Writing – review & editing. **Javad Abolfazli Esfahani:** Conceptualization, Writing – review & editing. **Nader Karimi:** Conceptualization, Writing – review & editing.



Fig. 14. Dependency of HNO mass fraction to enhancement of NH<sub>3</sub> in fuel blend in axial locations of L1 to L6 (x- axis is normalized with inlet fuel nozzle diameter).



Fig. 15. Radial distribution of H and O mass fractions in axial locations of L1 to L6 (x- axis is normalized with inlet fuel nozzle diameter).

# **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

# Data availability

No data was used for the research described in the article.

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