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10 **unsaturated finite slopes**

11

12 AUTHORS:

13 Thi Minh Hue Le<sup>1</sup>, Marcelo Sanchez<sup>2</sup>, Domenico Gallipoli<sup>3</sup>, Simon Wheeler<sup>4</sup>

14

15 AFFILIATIONS:

16 <sup>1</sup> Norwegian Geotechnical Institute. Sognveien 72, Oslo, Norway; Tel: +47 (0) 9300  
17 1834; Email: thi.le@ngi.no

18 <sup>2</sup> Zachry Department of Civil Engineering, Texas A&M University, 3136 TAMU,  
19 College Station, TX 77843-3136, United States; Tel. (+1) -979- 862 6604; Fax: (+1) -  
20 979- 862 7696; Email: [msanchez@civil.tamu.edu](mailto:msanchez@civil.tamu.edu)

21 <sup>3</sup> Laboratoire SIAME, Fédération IPRA, Université de Pau et des Pays de l'Adour, Allée  
22 du Parc Montaury, Anglet, 64600, France; Tel : +33 (0)559574280; Email:  
23 domenico.gallipoli@univ-pau.fr

24 <sup>4</sup> Infrastructure & Environment Research Division, School of Engineering, University of  
25 Glasgow, G12 8LT, Glasgow, United Kingdom; Tel: +44 (0)141 330 5201; Fax: +44  
26 (0)141 330 4557; Email: [Simon.Wheeler@glasgow.ac.uk](mailto:Simon.Wheeler@glasgow.ac.uk)

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47 **Abstract**

48

49 Water infiltration destabilises unsaturated soil slopes by reducing matric suction, which  
50 produces a decrease of material cohesion. If the porosity of the soil is spatially  
51 heterogeneous, a degree of uncertainty is added to the problem as water tends to follow  
52 preferential paths and produces an irregular spatial distribution of suction. This study  
53 employs the finite element method together with Monte Carlo simulations to quantify the  
54 effect of random porosity on the uncertainty of both the factor of safety and failure size of  
55 an unsaturated finite slope during and after a rainfall event. The random porosity is  
56 modelled using a univariate random field. Results show that, under partially saturated  
57 conditions, the random heterogeneity leads to a complex statistical variation of both  
58 factor of safety and failure size during the rainfall event. At any given time, the  
59 uncertainty about failure size is directly linked to the uncertainty about the position of the  
60 wetting front generated by infiltration. Interestingly, the statistical mean of the failed area  
61 is smallest when the mean of the factor of safety is lowest. In other words, the slope  
62 becomes more likely to fail but the size of the failure mass tends to be limited.

63 The study also investigates the sensitivity of failure uncertainty to external hydraulic  
64 parameters (i.e. initial water table depth, rainfall intensity) and internal soil parameters  
65 (i.e. permeability and water retention characteristics). In general, the sensitivity increases  
66 when the effect of these parameters on the spatial variation of suction is stronger.

67

68

69

## 70 **1 Introduction**

71 Catastrophic failures of soil slopes caused by rainfall infiltration are relatively common  
72 but their triggering mechanisms are still poorly understood. This is particularly true in  
73 unsaturated slopes where the spatial variability of suction and degree of saturation  
74 induces an uneven distribution of permeability inside the soil mass. This also means that,  
75 unlike in saturated soils, the permeability of unsaturated soils does not remain constant  
76 during the rainfall. The high non-linearity of the constitutive equations linking the soil  
77 suction (or saturation) to permeability and the coupling between soil porosity and degree  
78 of saturation make the numerical solution of these problems very challenging.

79

80 Further complexities are introduced by the heterogeneity of porosity, which influences  
81 the infiltration pattern and hence the stability of the slope. In a heterogeneous slope,  
82 water will preferably infiltrate through paths connecting high permeability areas, which  
83 in turn produces a spatially irregular distribution of suction and saturation inside the soil  
84 mass (Le et al. 2012). Soil elements experiencing an earlier loss of suction will also  
85 undergo an earlier reduction of strength compared to other elements where suction  
86 changes are slower. At any given time, the likely slip surface will therefore tend to pass  
87 through these weaker elements, which may result in a lower safety factor compared to a  
88 homogenous slope.

89

90 A relatively large number of probabilistic studies have investigated the effect of material  
91 uncertainties on the safety of dry or saturated slopes. Many of them have employed the  
92 finite element method (FEM), which is particularly suited to the description of spatial  
93 heterogeneity, to analyse the effect of strength variability on slope safety (Hicks 2005;  
94 Griffiths and Fenton 2004). Other studies have instead employed the limit equilibrium  
95 method (LEM) because of its simplicity (Pathak et al. 2007; El-Ramly et al. 2005).  
96 Stochastic studies of slope instabilities in randomly heterogeneous slopes have relied on  
97 Monte Carlo simulations to handle complicated geometries and variability patterns  
98 without requiring over-simplified assumptions. Results from these simulations, and from

99 practical observations, have repeatedly indicated that material heterogeneity affects  
100 strongly the stability of soil slopes (Alonso 1976; Babu and Mukesh 2004; El-Ramly et  
101 al. 2005; Griffiths and Fenton 2004; Griffiths and Marquez 2007; Hicks and Onisiphorou  
102 2005; Hicks and Samy 2002; Hicks and Spencer 2010; Mostyn and Li 1993; Mostyn and  
103 Soo 1992; Sejnoha et al. 2007; Cho 2009; Fenton and Griffiths 2005; Griffiths et al.  
104 2015). The majority of stochastic studies adopted the Monte Carlo approach because of  
105 its conceptual simplicity and its capability to handle complicated geometry and variability  
106 patterns without requiring over-simplified assumptions. A number of works based on  
107 Monte Carlo simulation have yielded a full description of the shearing processes and the  
108 probability of failure or the reliability of fully saturated heterogeneous slopes (Griffiths  
109 and Fenton 2004; Griffiths and Marquez 2007; Hicks and Onisiphorou 2005; Hicks and  
110 Samy 2002).

111 There have been a number of studies investigating the influence of rainfall intensity,  
112 water table and permeability on the stability of saturated slope (e.g., Tsaparas et al.  
113 (2002)). The main findings from these works cannot be directly applied to unsaturated  
114 slopes, because the flow characteristics in unsaturated soils are different from the ones  
115 observed under saturated conditions. Past studies on unsaturated slope stability are mostly  
116 limited to homogeneous soil properties and were conducted using different approaches,  
117 including analytical solution, the LEM and the FEM. Griffiths and Lu (2005) and Lu and  
118 Godt (2008) suggested a formula based on suction stress that takes into account both, the  
119 soil characteristics and the infiltration rate. The suction stress was then used to  
120 analytically predict the stability of an infinite unsaturated slope in a steady seepage  
121 condition. Ng and Shi (1998) conducted a LEM parametric study to investigate the effect  
122 of various hydraulic parameters, amongst others: permeability, rainfall intensity,  
123 infiltration duration and boundary conditions. It was observed that soil permeability and  
124 rainfall characteristics (i.e. intensity and duration) could have significant influences on  
125 the stability of unsaturated slopes. Importantly, the factor of safety can reduce  
126 considerably with the relative differences in magnitude between the soil permeability and  
127 the rainfall intensity and it might also depend on permeability anisotropy.

128  
129

130 Few studies have also attempted to incorporate material uncertainties into a stochastic  
131 analysis of partly saturated slopes. Among these studies, some are limited to the analysis  
132 of infinite slopes with one-dimensional random variations of permeability (Santoso et al.  
133 2011; Dou et al. 2014; Cho and Lee 2001; Cho 2014; Xia et al. 2017). For example, Dou  
134 et al. (2014) employed a Green-Ampt infiltration model to obtain a closed form of the  
135 limit state function of an infinite slope. The Monte Carlo simulation method was then  
136 used to study the influence of saturated permeability on slope failure during rainfall. Xia  
137 et al. (2017) adopted a stochastic method to predict the risk of failure of an infinite  
138 unsaturated slope subjected to rainfall. They proposed an analytic solution and compared  
139 it against a Monte Carlo simulation.

140

141 Sensitivity analyses looking at the effect of different factors (e.g. slope angle, water table  
142 position, soil air entry value, dry density and specific density) on slope failure were also  
143 conducted. Zhang et al. (2005) developed a coupled hydro-mechanical finite element  
144 model to study the effect of the variability of different constitutive parameters. Zhang et  
145 al. (2014) also extended this model to the analysis of rainfall intensity-duration and  
146 suggested a framework for predicting time-dependent failure probability. Arnold and  
147 Hicks (2010) studied the effect of the random variability of friction angle, cohesion,  
148 porosity, saturated permeability and air entry suction on the stability of a finite  
149 unsaturated slope. Phoon et al. (2010) proposed a probabilistic model of normalised soil  
150 water retention curve (SWRC), whose shape and air entry value were modelled by a  
151 correlated lognormal vector. The study did not however take into account the variability  
152 of saturated permeability. Santoso et al. (2011) further developed the SWRC model  
153 proposed in Phoon et al. (2010) by incorporating the saturated water content as an  
154 additional random variable. The Kozeny-Carman equation was adopted to link the  
155 random saturated water content to the saturated permeability. This approach implies that  
156 the shape of the SWRC and the saturated permeability are independent from one another,  
157 while in the present study they are coupled through the porosity as described later.

158

159 A limited number of authors have also investigated the depth of the failure zone. Alonso  
160 and Lloret (1983) showed that the slope angle marking the transition from shallow to

161 deep failure increases with soil dryness. Hicks et al. (2008) presented a three-dimensional  
162 stochastic study of the size of the sliding area in saturated slopes. Santoso et al. (2011)  
163 demonstrated instead that shallow failure mechanisms in randomly heterogeneous infinite  
164 unsaturated slopes cannot be predicted using a homogeneous slope model. Finally, Le et  
165 al. (2015) evaluated the effect of the standard deviation and correlation length of random  
166 porosity on the size of the sliding area in an unsaturated slope.

167

168 Following upon earlier studies, the present work investigates the effect of external and  
169 internal factors on the uncertainty of the factor of safety and failure size in unsaturated  
170 slopes with randomly heterogeneous porosity. These factors include external  
171 environmental conditions (i.e. water table depth and rainfall intensity) and internal soil  
172 parameters (i.e. saturated permeability and water retention characteristics). Importantly,  
173 unlike random saturated soils, preferential water pathways do not necessarily coincide  
174 with the most porous regions (Le et al. 2015). These regions might in fact exhibit smaller  
175 values of permeability because of lower saturation levels. A fully coupled hydro-  
176 mechanical FE code is adapted to perform the numerical simulations involving a finite  
177 slope. The Monte Carlo method is adopted to conduct the probabilistic study.

178

## 179 **2 Method**

### 180 **2.1 Model geometry**

181

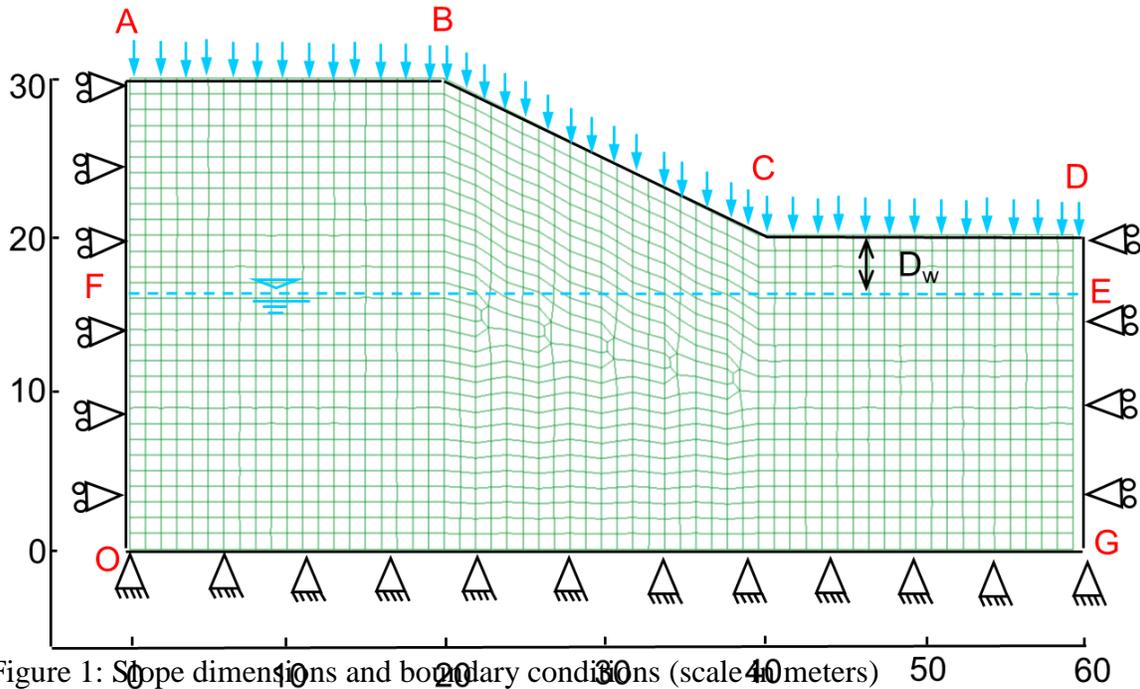


Figure 1: Slope dimensions and boundary conditions (scale 40 meters)

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186  
187  
188 The numerical model adopted in the present analysis consists of a slope with a 2:1  
189 gradient discretized into a finite element mesh of 1515 quadrilateral elements with four  
190 integration point and an average area of  $\sim 1\text{m}^2$  (Figure 1). The finite element  
191 CODE\_BRIGHT software (Olivella et al. 1996; UPC 2010) was adopted to conduct the  
192 numerical analyses. This software allows fully coupled thermo-hydro-mechanical  
193 simulations of boundary value problems in unsaturated soils. Thermal processes are  
194 however not considered in this study, which focuses exclusively on coupled hydro-  
195 mechanical processes.

196

197 A mesh sensitivity analysis was initially performed under saturated conditions, which  
198 confirmed the accurate estimation of the safety factor by the model shown in Figure 1 (Le  
199 2011). The suitability of the mesh was further verified in unsaturated conditions against  
200 commercial software (GEO-SLOPE International Ltd) using the limit equilibrium  
201 method. For a given rainfall, the commercial software produced similar changes of the  
202 factor of safety compared to the adopted finite element model (Le et al. 2015).

203

## 204 2.2 Hydraulic and mechanical models

205 The hydraulic constitutive models adopted in this study are presented in Eqs. 1 to 5:

$$206 \quad S_e = \frac{S - S_r}{S_s - S_r} = \left( 1 + \left( \frac{s}{s_e} \right)^{1-m} \right)^{-m} \quad (1)$$

$$207 \quad s_e = s_{eo} \exp(\eta(\phi_o - \phi)) \quad (2)$$

$$208 \quad k_s = k_{so} \frac{\phi^3}{(1-\phi)^2} \frac{(1-\phi_o)^2}{\phi_o^3} \quad (3)$$

$$209 \quad k_r = \sqrt{S_e} (1 - (1 - S_e^{1/m})^m)^2 \quad (4)$$

$$210 \quad \mathbf{q} = -k_s k_r \left( \frac{u_w}{\rho_w g} + z \right) = -k_u \nabla \left( \frac{u_w}{\rho_w g} + z \right) \quad (5)$$

211

212

213 This work employs the van Genuchten (1980) model for the soil water retention curve  
 214 (SWRC) (Eq. 1-2), the Kozeny's relationship (Kozeny 1927) between saturated  
 215 permeability and porosity (Eq. 3) and the van Genuchten and Nielsen (1985) model for  
 216 the unsaturated relative permeability (Eq. 4). The unsaturated permeability  $k_u$  is then the  
 217 product of the saturated and relative permeabilities (i.e.  $k_u = k_s k_r$ ) while the unsaturated  
 218 flow  $\mathbf{q}$  is calculated using the generalised Darcy's law (Eq. 5). The above models can  
 219 realistically describe unsaturated flow in a simple and numerically stable way, which is  
 220 highly desirable when dealing with finite element simulations. Nevertheless, they rely on  
 221 the simplifying assumption that capillarity dominates the hydraulic regime and that other  
 222 forces linked to adsorptive phenomena are negligible.

223

224 The SWRC (Eq. 1) relates the effective degree of saturation  $S_e$  to suction  $s$  through the air  
 225 entry suction parameter  $s_e$  and the retention gradient  $m$  (van Genuchten 1980). The value  
 226 of  $S_e$  is calculated as a function of the current degree of saturation  $S$ , the maximum degree  
 227 of saturation  $S_s$ , and the residual degree of saturation  $S_r$ . **The effect of heterogeneity is**  
 228 **introduced by relating the parameter  $s_e$  to porosity  $\phi$  through the parameter  $\eta$  (Eq. 2) that**

229 controls the rate at which  $s_e$  deviates from its reference value  $s_{e0}$  when  $\phi$  deviates from its  
 230 reference value  $\phi_0$  (Rodríguez et al. 2007; Zandarín et al. 2009). Similarly, Kozeny's  
 231 equation (Eq. 3) describes the deviation of the saturated permeability  $k_s$  from its reference  
 232 value  $k_{s0}$  when  $\phi$  deviates from its reference value  $\phi_0$  (Kozeny 1927). The van Genuchten  
 233 and Nielsen (1985) permeability curve (Eq. 4) relates instead the relative permeability  $k_r$   
 234 to the effective degree of saturation  $S_e$ , and therefore indirectly to porosity  $\phi$ , through the  
 235 gradient  $m$  of the SWRC curve. The symbols  $u_w$ ,  $\rho_w$ ,  $g$  and  $z$  indicate the pore water  
 236 pressure, the water density, the gravitational acceleration and the elevation coordinate,  
 237 respectively. The water retention behaviour and permeability are therefore spatially  
 238 heterogeneous which influences the hydraulic processes within the soil masses . More  
 239 details about these relationships can be found in UPC (2010).

240

241 Unless otherwise stated, the base values of  $m$ ,  $k_{s0}$ ,  $s_{e0}$  and  $\eta$  are constant and equal to the  
 242 values shown in Table 1. These values are about the middle of their respective typical  
 243 range of variation (i.e. those values that are physically possible and are of interest in  
 244 practically applications) to avoid unrepresentative results (Bear 1972; van Genuchten  
 245 1980; Zandarín et al. 2009). The base value of  $k_{s0}=10^{-5}$  m/s lies in the upper permeability  
 246 range of layered clays or clayey silts. The choice of a relatively high  $k_{s0}$  facilitates  
 247 numerical simulations by easing the steep change of pore pressure across the wetting  
 248 front. During the sensitivity analysis, the parameters  $k_{s0}$ ,  $\eta$  and  $m$  are varied in their  
 249 typical range to investigate the effect on slope stability. In Eq. 1, the values of  $S_s$  and  $S_r$   
 250 are equal to 1 and 0.01, respectively.

251

252 A linear elastic model with an extended Mohr-Coulomb (MC) failure criterion (Eq. 6) is  
 253 adopted to simulate the mechanical behaviour of the unsaturated soil (Fredlund et al.  
 254 1978):

255

$$256 \quad \tau = c' + \sigma \tan \phi' + s \tan \phi^b \quad (6)$$

257

258 Eq. 6 reflects the dependency of the shear stress at failure  $\tau$  on net normal stress  $\sigma$  and  
 259 suction  $s$  through the effective friction angle  $\phi'$ , effective cohesion  $c'$  and a parameter  
 260 controlling the increase in shear strength with suction  $\phi^b$ . The cohesive component of  
 261 strength provided by suction (i.e. the 3<sup>rd</sup> term in Eq. 6) reduces with decreasing  $s$  and  
 262 becomes zero for a fully saturated soil (i.e.  $s = 0$ ). In reality, the value of  $\phi^b$  has been  
 263 shown experimentally not to be constant but to decrease with increasing  $s$  (Escario and  
 264 Saez 1986; Gan et al. 1988) starting from  $\phi'$  in saturated conditions. In particular, Gan et  
 265 al. (1988) suggested that, as the soil desaturates, the value of  $\phi^b$  decreases up to a  
 266 relatively constant value. For simplicity, however, this study assumes a constant value of  
 267  $\phi^b$ .

268

269 The assumed values of  $c'$ ,  $\phi'$  and  $\phi^b$  are typical of clays and are based on those reported  
 270 by Bishop et al. (1960) for boulder clay and by Gan et al. (1988) for a compacted glacial  
 271 till. The elastic parameters (i.e. Young's modulus  $E$  and Poisson's ratio  $\nu$ ), are also  
 272 related to typical values observed in clayey soils, and chosen within their respective  
 273 ranges (Zhu 2014). The variation of porosity may also influence mechanical behaviour,  
 274 but this aspect is not considered in this study. **The mechanical parameters are therefore**  
 275 **assumed to be homogeneous (spatially uniform) and are set equal to the values listed in**  
 276 **Table 1.** This assumption facilitates the investigation of the effect of porosity  
 277 heterogeneity on the hydraulic behaviour by isolating it from other effects.

278

279 A non-associated flow rule with zero dilatancy is assumed, which means that no plastic  
 280 volumetric strains occur during yielding. Moreover, a viscoplastic integration algorithm  
 281 is used to update the stress field during plastic loading (Olivella et al. 1996).

282

283 Table.1: Base values of soil parameters adopted in the numerical analyses

Hydraulic model			Mechanical model		
Symbol	Units	Value	Symbol	Units	Value
$m$		0.2	$E$	kPa x 10 <sup>3</sup>	100

$\eta$		5	$\nu$		0.3
$\phi_o$		0.333	$\phi'$	°	20
$k_{so}$	m/s	$10^{-5}$	$c'$	kPa	5
$s_{eo}$	kPa	20	$\phi^b$	°	18

284

285 As shown in Eq. 6,  $\tan\phi^b$  controls the increase in shear stress at failure with suction,  
 286 which provides an additional source of cohesive strength with respect to the effective  
 287 cohesion  $c'$ . Therefore, when implementing the shear strength reduction technique for  
 288 estimating the factor of safety ( $FoS$ ), the same reduction is applied to all strength  
 289 parameters ( $c'_{actual}$ ,  $\tan\phi'_{actual}$ ,  $\tan\phi^b_{actual}$ ) to obtain the corresponding values at failure  
 290 ( $c'_{fail}$ ,  $\tan\phi'_{fail}$ ,  $\tan\phi^b_{fail}$ ) according to the following definition of  $FoS$  for unsaturated soils:

291

$$292 \quad FoS = \frac{c'_{actual}}{c'_{fail}} = \frac{\tan\phi'_{actual}}{\tan\phi'_{fail}} = \frac{\tan\phi^b_{actual}}{\tan\phi^b_{fail}} \quad (7)$$

293

294 The use of Eq. 7 in conjunction with the *FE* program CODE\_BRIGHT has been verified  
 295 against the Limit Equilibrium Method by using the commercial software SeepW and  
 296 SlopeW (GEO-SLOPE International Ltd) and has been shown to produce comparable  
 297 values of  $FoS$  (Le 2011, Le et al. 2015). More details about the application of the shear  
 298 strength reduction method using CODE\_BRIGHT can be found in Le (2011) and Le et al.  
 299 (2015).

300

### 301 **2.3 Boundary conditions and simulation process**

302 At the very start of the analysis, gravity is applied to an initially weightless slope to  
 303 establish the initial stress distribution due to self-weight. The acceleration of gravity is  
 304 increased from zero to the standard value of  $9.8 \text{ m/s}^2$  over a 'fictitious' time (UPC 2010).  
 305 The random porosity field is introduced prior to applying gravity, so that the initial stress

306 distribution takes into account the variation of the soil unit weight due to material  
307 heterogeneity.

308

309 The initial distribution of pore water pressure  $p_w$  is assumed hydrostatic in equilibrium  
310 with the water table. The water table is fixed at 5 m below the slope toe, except for those  
311 analyses where the effect of water table depth is investigated. The pore air pressure is  
312 assumed constant and equal to the atmospheric pressure (i.e.  $p_a=0$ ) and the suction  $s$  is  
313 therefore equal to the negative value of the pore water pressure (i.e.  $s=-p_w$ ). The initial  
314 suction is therefore largest at the crest of the slope  $AB$  and equal to  $s_{max}=150$  kPa under  
315 hydrostatic conditions. This level of surface suction is typically encountered in arid or  
316 semi-arid countries such as Australia (e.g., Cameron et al. (2006)). The assumption of an  
317 initially hydrostatic pore pressure distribution ignores the potential presence of  
318 evaporation at ground level. This simplification is acceptable in the context of this work,  
319 whose objective is to analyse the sensitivity of the stability of unsaturated slopes to  
320 different parameters rather than describing the hydrological and failure regimes of a real  
321 case.

322

323 A rainfall of constant intensity is then applied at the boundary  $ABCD$  over 10 days  
324 (Figure 1). This boundary condition imposes a constant rate of infiltration into the soil as  
325 long as the pore water pressure at the boundary is negative (i.e. as long as suction is  
326 positive). If the pore water pressure becomes equal or larger than zero, the boundary  
327 condition shifts to a constant zero pore water pressure to avoid the build-up of a hydraulic  
328 head at the ground surface. This type of boundary condition is often referred to as a  
329 “seepage” boundary condition and is further described in CODE\_BRIGHT Users’  
330 Manual (UPC 2010) or Le et al. (2012). After 10 days, the rainfall is stopped and the  
331 boundary  $ABCD$  is assumed impermeable but the simulation is continued for another 355  
332 days to allow the redistribution of pore water pressure back to a hydrostatic condition.  
333 The boundaries  $OA$ ,  $OG$  and  $GD$  are assumed impermeable during and after the rainfall,  
334 which causes the infiltrated water to accumulate inside the soil domain and the water  
335 table to rise. This describes a situation in natural slopes where surrounding soils have low  
336 permeability or neighbouring areas have poor drainage capacity (e.g., due to a blocked

337 drain). Such a condition can indeed be critical for slope stability in reality. If evaporation  
338 and/or dissipation were allowed, the water table position would be affected depending on  
339 the considered assumptions. For example, if high rates of evaporation are assumed the  
340 rise of the water table will be strongly affected, leading to an eventual little water  
341 accumulation in the slope domain and therefore to a practically stable position of the  
342 water table during the rainfall. Then, the changes of the safety factor and size of failure  
343 mass during the rainfall would be less than the results obtained in this study. In addition,  
344 the values of these parameters after the rainfall would be almost the same as at the  
345 beginning of the rainfall. Similar reasoning can be used with respect to the inclusion of  
346 dissipation in the simulations. The mechanical boundary conditions are also indicated in  
347 Figure 1.

348

349 The Monte Carlo analysis involves the generation of multiple random porosity fields that  
350 are mapped onto the *FE* mesh shown in Figure 1. These *FE* meshes with different  
351 random porosity fields constitute the “realisations” of the Monte Carlo analysis. Each  
352 realisation is analysed in two consecutive stages corresponding to: *i*) the calculation of  
353 the pore water pressure and stress fields at distinct times during or after the rainfall; and  
354 *ii*) the application of the shear strength reduction technique (*SRT*) to the calculated pore  
355 water pressure and stress fields to determine the factor of safety (*FoS*) and sliding area  
356 ( $A_s$ ) at a given time.

357

358 Note that, in stage *i*), soil deformations are fully coupled with pore water flow and the  
359 equations of equilibrium and hydraulic continuity are solved simultaneously in  
360 CODE\_BRIGHT. The nonlinear equations associated with flow and mechanical  
361 problems are solved in a fully coupled manner using the New-Raphson method (Olivella  
362 et al., 1996). This implies that as the rainfall seeps into the unsaturated soil, suction  
363 (and/or positive pore water pressure) changes will induce net (or effective) stresses  
364 changes. This in turn induces deformations in the soil elements. These deformations  
365 cause changes in the soil porosity, which lead to changes in intrinsic permeability and air  
366 entry value through equations 2 and 3, respectively. The new permeability and air

367 entry value influence the water flows through equation 1, 4 and 5. The  
368 unsaturated/saturated flow and the mechanical deformations are therefore truly coupled.

369

370 Eight points in time are selected to extract the corresponding fields of stresses and pore  
371 water pressure to be used in the subsequent shear strength reduction stage. These include  
372 four times during the rainfall (i.e. 0, 0.5, 5, 10 days) and four times after the rainfall (i.e.  
373 15, 20, 100 and 365 days). The selected times aim at capturing the changes in the failure  
374 mechanism associated with a significant variation of the pore water pressure  $p_w$  field.

375

376 Note that the *SRT* analysis is simply a numerical technique used in stage *ii*) to estimate  
377 the factor of safety  $FoS$  and sliding area  $A_s$  corresponding to the field of stresses and pore  
378 water pressures calculated at a given time. During a *SRT* analysis, the calculated pore  
379 water pressures field is fixed at every mesh node while the calculated stresses and strains  
380 fields are imposed as initial conditions. The shear strength parameters are then reduced  
381 by a factor that is initially equal to one and subsequently augmented in steps of 0.01 until  
382 failure. Failure corresponds to the detection of significant movements on the slope  
383 surface. The value of the reduction factor at this point is assumed to coincide with the  
384  $FoS$  of the slope (Eq. 7). Note that the above methodology allows the natural  
385 development of the slip surface through the weakest path within the soil domain, which is  
386 an advantage compared with limit equilibrium methods where the shape of the slip  
387 surface is instead assumed. Le et al. (2015) provided detailed explanation of the criteria  
388 used to detect the failure mechanism.

389

390 After failure, the number of mesh nodes that have moved substantially is counted to  
391 compute the sliding area of the slope (Le et al. 2015). One node corresponds to a region  
392 that is the sum of one quarter of each of the four elements sharing that node. Since the  
393 mesh mostly consists of square or parallelogram elements of  $1 \text{ m}^2$  (Figure 1), the area  
394 allocated to each node is approximately  $1 \text{ m}^2$  and the number of “failed” nodes provides a  
395 reasonably good estimation of the sliding area  $A_s$  in  $\text{m}^2$ . This is clearly an approximation  
396 because the nodes on the boundary of the failed region contribute less area than the inner  
397 nodes. Nevertheless, this approximation is considered acceptable as the present study

398 focuses on a sensitivity analysis rather than on the accurate determination of the sliding  
399 area. For real slopes, it is recommended that  $A_s$  is estimated more accurately either by  
400 using a finer mesh or by directly measuring the area of the failed region.

### 401 **3 Random porosity field**

402 Porosity  $\phi$  is probably one of the most easily measured soil parameters exhibiting spatial  
403 variability (Le et al. 2013). Porosity values are theoretically bounded between 0 and 1,  
404 thus they should be represented by a bounded random distribution such as the tanh-  
405 bounded function. This distribution requires 4 parameters which are a lower bound, an  
406 upper bound, the location parameter (equal to 0 when random variable is symmetric  
407 about the midpoint of the variable range) and a scale parameter which increases with  
408 increasing level of variability. The bounded distributions are mathematically complex so  
409 a different approach is employed in the present work by generating an univariate random  
410 field of void ratio  $e$  instead of porosity  $\phi$ . The void ratio can take any positive value and  
411 may thus be modelled by a log-normal probability function (Baecher and Christian 2003;  
412 Lacasse and Nadim 1996). The generated random field of void ratio is then converted  
413 back into a random field of porosity by using the relationship  $\phi=e/(1+e)$ . This equation  
414 implies that the random field does not generate any value of porosity equal to zero. Such  
415 a value is considered unrealistic for the size of the mesh considered in this study.

416

417 Random fields of void ratio are produced by using the Local Average Subdivision (LAS)  
418 algorithm and the Markov auto-correlation function (Fenton 1990). The Local Average  
419 Subdivision (LAS) method (Fenton, 1990) involves a recursive subdivision process. The  
420 original domain is first subdivided into equal sized area, then each area is divided again  
421 into smaller areas and this process keeps going until the desirable resolution is achieved.  
422 At every stage of subdivision, random values are generated for each area with the  
423 variance and covariance structure inherently related to the size of the subdivided area  
424 relative to the original domain. Both the LAS algorithm and the Markov function have  
425 already been used in geotechnical engineering (Fenton 1990; Griffiths and Fenton 2004).  
426 The random field is generated over a regular grid covering a rectangular area with

427 dimensions equal to the largest width and height of the soil domain. The grid is then  
428 superimposed on the finite element mesh, so that the bottom left corners of the grid and  
429 mesh coincide. An algorithm is subsequently executed to identify the cell in the random  
430 field grid with the closest centroid to the centroid of each finite element. The void ratio of  
431 the finite element is then taken to coincide with the random value of that cell. Le (2011)  
432 explains in detail the procedure to verify that statistical parameters are correctly  
433 transferred in the above mapping process.

434

435 The effect of the statistical parameters governing the random distribution of void ratio  $e$   
436 (i.e. mean  $\mu(e)$ , standard deviation  $\sigma(e)$  and correlation length  $\theta(e)$ ) were studied in detail  
437 in Le et al. (2015). In this study, the values of the mean  $\mu(e)$ , coefficient of variation  
438  $COV_e = \sigma(e)/\mu(e)$  and correlation length  $\theta(e)$  are therefore kept constant and equal to 0.5,  
439 0.8 and 8 m, respectively (which correspond to  $\mu(\phi) = 0.3$ ,  $COV_\phi = 0.46$  and  $\theta(\phi) = 8$  m).  
440 The effect of  $COV_e$  and  $\theta(e)$  has been investigated in another study (Le et al. 2015). The  
441 chosen values for  $COV_e$  and  $\theta(e)$  aim to avoid too large or too small effect of these  
442 parameter on the results, and increase the possibility of observing the effect of porosity  
443 heterogeneity on suction distribution within the slope.

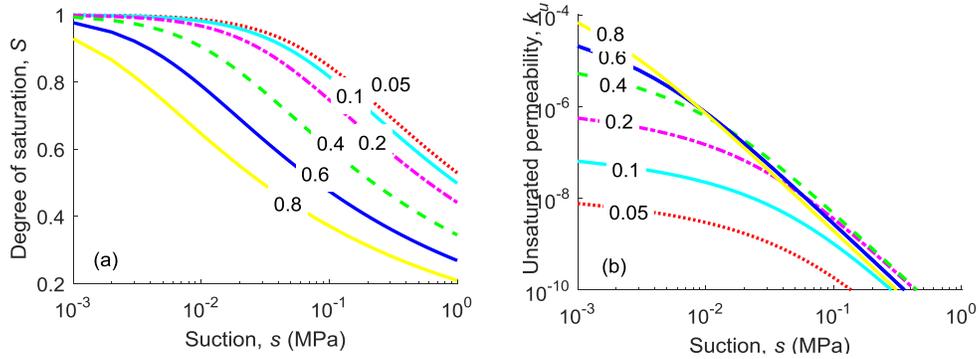
444

445 Figures 2a and 2b show the influence of porosity on the SWRC and  $k_u$  curves calculated  
446 using Eqs. 1-4 and the input parameters are listed in Table 1. Six values of porosity, from  
447 0.05 to 0.8, are considered. A value of porosity outside this range is quite unlikely  
448 considering the coefficient of variation adopted in this study. Based on Figure 2a, the  
449 initial degree of saturation near the crest of the slope (i.e.  $s \approx 150$  kPa) varies between 0.3  
450 and 0.8 with a corresponding value of  $k_u$  in the range  $10^{-10}$ – $10^{-9}$  m/s.

451

452 A heterogeneous porosity field therefore generates non-uniform distributions of degree of  
453 saturation and permeability (in addition to a non-uniform distribution of specific weight),  
454 which leads to an irregular advancement of the wetting front and an uneven distribution  
455 of pore water pressures. This affects the distribution of shear strength, which is controlled  
456 by pore water pressure (in addition to the distribution of stresses, which is governed by

457 the overburden weight) and has an impact on the factor of safety of the slope as well as  
 458 on the size of the sliding mass.  
 459



460  
 461 Figure 2. Influence of porosity on the variation of degree of saturation (a) and unsaturated  
 462 permeability (b) with suction.

463 Noticeably, the degree of saturation (Figure 2a) decreases with increasing porosity while  
 464 the unsaturated permeability (Figure 2b) increases with increasing porosity. The latter  
 465 (i.e.  $k_u$ ) is however little affected when suction is above 20 kPa and the porosity is higher  
 466 than 0.2. This implies that, in unsaturated soils, the higher porosity regions are not  
 467 necessarily the most permeable ones, as it is instead the case in saturated soils.  
 468

## 469 4 Influence of hydraulic characteristics

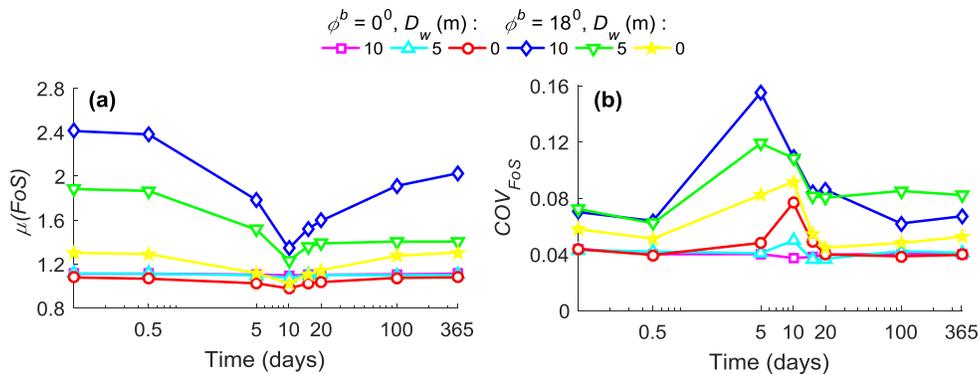
### 470 4.1 Water table depth

471 The initial suction of the soil affects both its degree of saturation and unsaturated  
 472 permeability (Eqs. 1 and 4), which makes the initial position of the water level ( $D_w$ ) an  
 473 important factor to consider. Three values of water table depth measured with respect to  
 474 the toe of the slope are investigated in this section, namely 0, 5 and 10 m. Under  
 475 hydrostatic conditions, these depths correspond to the three maximum values of initial  
 476 suction at the crest of the slope of 100, 150 and 200 kPa, respectively. For each depth,  
 477 two analyses are compared: one considering the effect of suction on shear strength, i.e.  
 478  $\phi=18^\circ$ , and one neglecting this effect, i.e.  $\phi=0$ .

479 The evolution of the mean and coefficient of variation of  $FoS$ , i.e.  $\mu(FoS)$  and  $COV_{FoS}$ ,  
 480 are presented in Figures 3a and 3b, respectively. When the effect of suction is considered  
 481 (i.e.  $\phi^b=18^\circ$ ), the  $\mu(FoS)$  progressively decreases during the rainfall, because of the  
 482 reduction in shear strength triggered by the reduction of suction in the unsaturated region  
 483 but also because of the build-up of positive pore water pressures in the saturated area at  
 484 the slope toe. In all the analyses, the lowest value of  $\mu(FoS)$  occurs just before the end of  
 485 the rainfall. The  $\mu(FoS)$  then recovers over the post-infiltration period (i.e. day 10 to 365),  
 486 because of the suction increase caused by the downward drainage and the consequent  
 487 dissipation of positive pore water pressure. The final  $\mu(FoS)$  values (i.e. at day 365) are  
 488 lower than the initial ones because of the rise of water table induced by the accumulation  
 489 of infiltrated water.

490

491 For the case of  $\phi^b=18^\circ$ , the  $\mu(FoS)$  consistently increases with increasing  $D_w$  because of  
 492 the increase in shear strength with growing suction. As rainfall progresses, the slope with  
 493 the deepest initial water table (i.e.  $D_w=10$  m) loses the largest amount of suction, leading  
 494 to the most substantial reduction in  $\mu(FoS)$  from about 2.4 to 1.3 over the 10 days of the  
 495 rainfall. Instead, the  $\mu(FoS)$  of the slope with the shallowest initial water table (i.e.  $D_w=0$   
 496 m) reduces much less from about 1.3 to 1.0 over the same time.

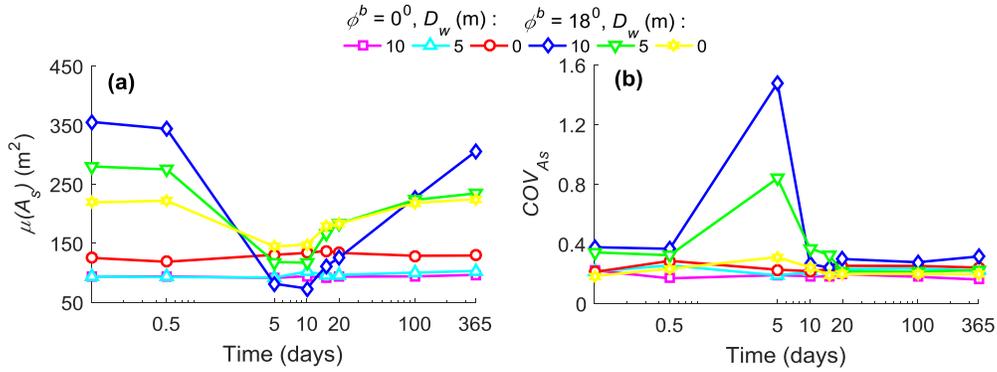


497

498 Figure 3: Time evolution of  $FoS$  in terms of mean (a) and coefficient of variation (b).

499 Analyses: influence of water table depth  $D_w$ .

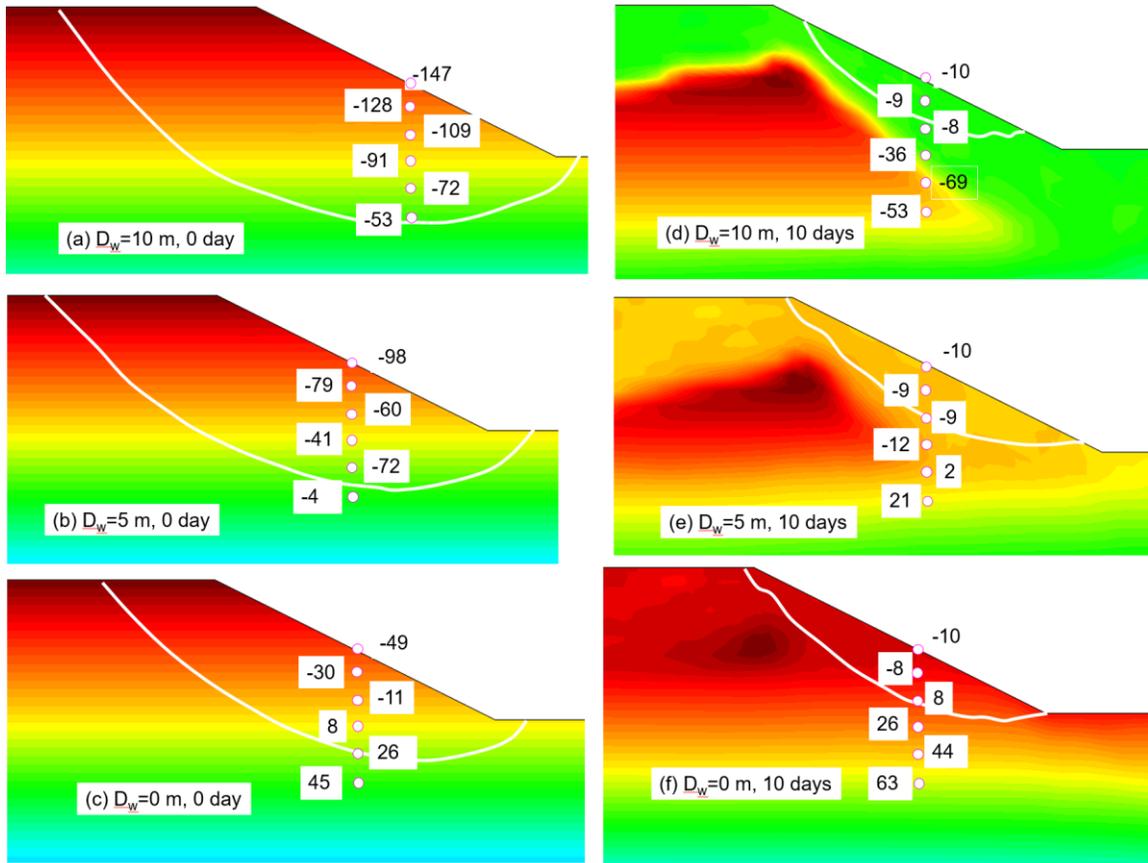
500



501  
 502 Figure 4. Time evolution of  $A_s$  in terms of mean (a) and coefficient of variation (b).  
 503 Analyses: influence of water table depth  $D_w$ .

504 Similar results are shown in Figures 4a and 4b but in terms of  $\mu(A_s)$  and  $COV_{A_s}$ ,  
 505 respectively. When the effect of suction is included ( $\phi^b=18^\circ$ ), the value of  $\mu(A_s)$   
 506 consistently decreases during the rainfall (though at different rates depending on the  $D_w$   
 507 value) and reaches a plateau between 5 and 10 days before increasing again during the  
 508 post-infiltration period. The reason behind this behaviour is that, at the start of the  
 509 rainfall, the shallow soil region exhibits considerable strength arising from the high  
 510 suction, which 'pushes' the slip surface to deeper layers in the search of a 'weak' path  
 511 (Figure 5). However, after a rainfall time between 5 and 10 days, the shallow soil  
 512 experiences a dramatic loss of suction and therefore becomes significantly weaker than  
 513 the deeper soil. This in turn promotes the formation of a slip surface through the wetted  
 514 shallow soil layer, which explains why  $A_s$  tends to decrease (Figure 5b, 5d, 5f).

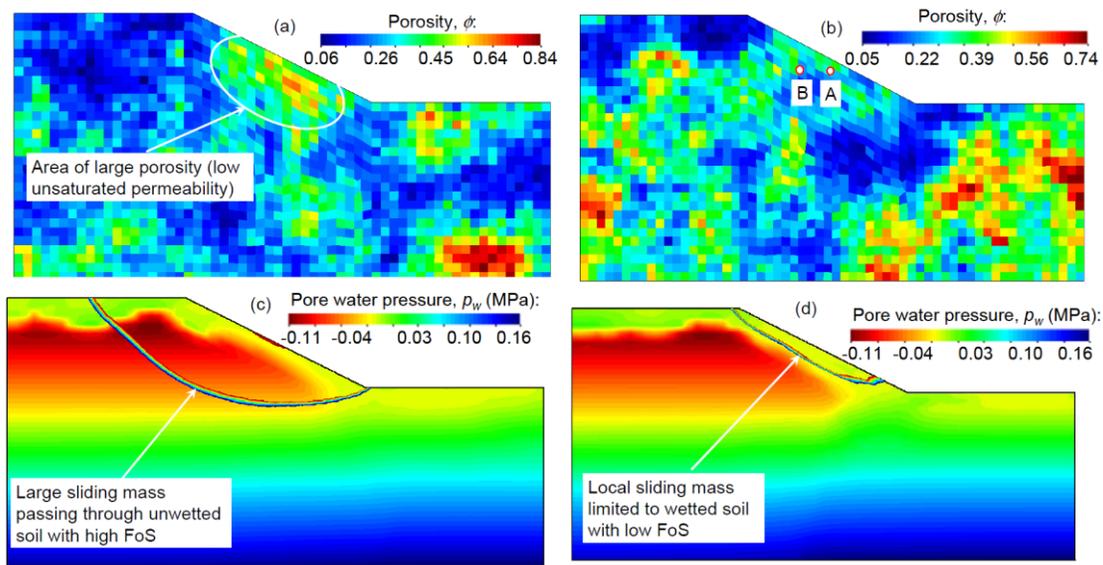
515  
 516 For the case of  $\phi^b=18^\circ$ , the values of  $\mu(A_s)$  are higher for larger values of  $D_w$ , both at the  
 517 beginning (i.e. 0 to 0.5 day) and at the end (i.e. 100 to 365 days) of the analysis, because  
 518 of the larger soil suction associated to a depressed water table (Figure 5a, 5c and 5e).  
 519 During the course of the rainfall, the wetted area decreases in depth with increasing  $D_w$   
 520 because of the higher initial suction, and hence the lower degree of saturation and  
 521 permeability, which delays water infiltration (Figure 5b, d and f). This explains the higher  
 522 value of  $\mu(A_s)$  with smaller  $D_w$  between 5 and 10 days (Figure 4a).



523

524 Figure 5. Contour maps of  $p_w$  and slip surfaces for different  $D_w$  at different times  
 525 ( $\phi=18^\circ$ ). The  $p_w$  values shown in labels are in kPa. The  $p_w$  colour scale is not the same  
 526 for all contour plots.

527



528

529 Figure 6: Porosity distributions of sample realisations with significantly different failure  
 530 mechanisms (a, b) and contour maps of  $p_w$  with sliding surfaces at 5 days (c, d). Results  
 531 correspond to  $\phi^b=18^\circ$  and  $D_w=5$  m.

532 For the case of  $\phi^b=18^\circ$ , the sliding area at 5 days varies over a wide range of values  
 533 depending on the depth of the wetting front in each realisation. There appears to be a  
 534 'critical' depth such that, when the wetting front moves below it, the sliding area is  
 535 confined to the superficial wetted region (Figure 6b and 6d). In this case, the  $FoS$  tends to  
 536 be low, because the suction of the 'wetted' elements is relatively low (Figure 56d).  
 537 Conversely, if the wetting front is shallower than the 'critical' depth, the slip surface tends  
 538 to be deep seated (Figure 6c), like at the start of the rainfall, with a large  $FoS$  due to the  
 539 high suction along the slip surface. This case might correspond to the existence of a low  
 540 permeability layer that prevents the advancement of the wetting front (Figure 6a). The  
 541 equal occurrence of both these two extremes (i.e. shallow versus deep slip surfaces)  
 542 causes the large values of  $COV_{FoS}$  and  $COV_{As}$  at 5 days. At 10 days, the wetting front is  
 543 likely to have passed the 'critical' depth and hence the majority of slip surfaces is  
 544 confined to the superficial wetted region, which explains the consistent decrease in  
 545  $COV_{FoS}$  and  $COV_{As}$ . An exception to this behaviour is the  $COV_{FoS}$  for the case of  $D_w=0$  m,

546 which peaks at 10 days because of the dominant destabilizing effect of positive pore  
547 pressure build-up at the slope toe.

548 The peak values of  $COV_{FoS}$  and  $COV_{A_s}$  significantly increase with increasing  $D_w$  implying  
549 that the factor of safety and the size of the sliding area become more variable between  
550 realisations. After the peak, the values of  $COV_{FoS}$  and  $COV_{A_s}$  decrease because of water  
551 drainage causing an increase of suction in the unsaturated region and a dissipation of  
552 positive pore pressures in the saturated region, which reduce the difference between  
553 realisations.

554

555 When the effect of suction on shear strength is not considered (i.e.,  $\phi^b=0$ ), Figure 3 shows  
556 that the  $\mu(FoS)$  is virtually constant for all three  $D_w$  values, with only a slight decrease at  
557 day 10 for  $D_w=0$ , while the  $COV_{FoS}$  increases slightly with decreasing  $D_w$  between 5 and  
558 20 days. The build-up of positive pore water pressures with decreasing  $D_w$  is the main  
559 reason behind this trend given that a larger portion of the slip surface passes through the  
560 saturated region as the initial water table is shallower. Figure 4 shows that  $\mu(A_s)$  and  
561  $COV_{A_s}$  remain fairly constant over time. Inspection of displacement contours (not shown  
562 here) reveal that the sliding areas are very similar for  $D_w=5$  m and  $D_w=10$  m and do not  
563 practically change over time. When the water table is at the ground surface, sliding areas  
564 tend to be slightly larger due to the additional stabilizing effects provided by the weight  
565 of water in the saturated part of the slope.

566

567 Similar patterns of variation with time of the mean and coefficient of variation of both  
568  $FoS$  and  $A_s$  were observed in all cases hereafter, hence they will not be discussed further.  
569 The comments will instead focus on the sensitivity of the results to the parameters under  
570 study.

## 571 **4.2 Saturated permeability**

572 The reference saturated permeability  $k_{so}$  controls the infiltration rate and influences the  
573 advancement of the wetting front together with the distribution of pore water pressures. A  
574 range of realistic  $k_{so}$  values, from  $10^{-4}$  m/s (e.g. pervious well sorted sands) to  $10^{-7}$  m/s

575 (e.g. silts or layered clays), is investigated in this section to gain insights into the  
576 influence of this parameter on slope stability.

577

578 Figure 7 and Figure 8 show similar variations of  $\mu(FoS)$ ,  $\mu(A_s)$ ,  $COV_{FoS}$  and  $COV_{A_s}$  over  
579 time as observed in the previous section, except for the lowest value of the reference  
580 permeability (i.e.  $k_{so}=10^{-7}$  m/s). In this case, almost no water infiltrates the soil and all  
581 curves remain practically flat over the entire simulation period.

582

583 Notably, the variation of  $\mu(FoS)$  and  $COV_{FoS}$  with  $k_{so}$  is not monotonic (Figure 7) and the  
584 intermediate value of  $k_{so}$  (i.e.  $10^{-5}$  m/s) causes the largest average drop of factor of safety  
585 as well as the widest variability between realisations (i.e. lowest  $\mu(FoS)$  and highest  
586  $COV_{FoS}$  for the period 5 to 10 days). This is because the highest value of  $k_{so}$  (i.e.  $10^{-4}$  m/s)  
587 facilitates water flow leading to smaller gradients of pore pressure together with smaller  
588 drops in suction, which results in smaller reductions of shear strength. Conversely, the  
589 intermediate value of  $k_{so}$  (i.e.  $10^{-5}$  m/s) generates larger gradients of pore pressure with  
590 bigger suction drops, which allows the sliding surface to remain inside the wetted region  
591 at the surface. This explains the lower values of  $\mu(FoS)$  and  $\mu(A_s)$  and the higher values of  
592  $COV_{FoS}$  and  $COV_{A_s}$  for  $k_{so}=10^{-5}$  m/s compared to  $k_{so}=10^{-4}$  m/s. The evolution of pore  
593 water pressures at the two sampling points shown in Figure 9a confirms the larger suction  
594 drops at 10 days for  $k_{so}=10^{-5}$  m/s compared to  $k_{so}=10^{-4}$  m/s (Figure 10).

595

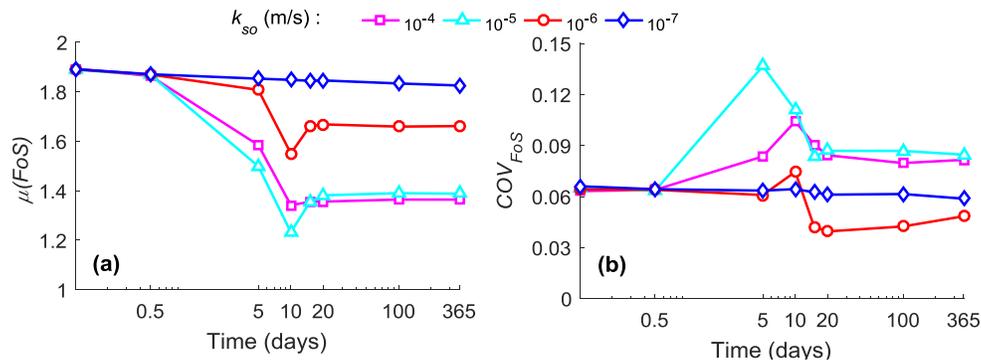
596 The lower value of  $k_{so}$  (i.e.  $10^{-6}$  m/s) limits infiltration and restricts the water movement  
597 to a very shallow layer along the slope face (Figure 9b). In this case, most of the suction  
598 loss is limited to the narrow top region (Figure 10a) while a wider wetted region develops  
599 at the slope toe (Figure 9b). Slip surfaces concentrate in this wetted region, which results  
600 in smaller values of  $COV_{FoS}$  with higher values of  $\mu(FoS)$  compared to the previous two  
601 cases (Figure 7). Moreover, the value of  $\mu(A_s)$  shows a sharp drop at 10 days because of  
602 the dominant failure mode cutting through the wetted region above the slope toe (Figure  
603 8a). The  $COV_{A_s}$  attains a sharp peak at 10 days (Figure 8b) because of the contrast  
604 between the majority of realisations predicting a small sliding area constrained to the  
605 wetted region and few others predicting a very large value of the sliding area. The latter

606 scenario is observed when the area near the slope toe is dominated by highly permeable  
 607 soil.

608

609 The drop of  $\mu(A_s)$  and the peak of  $COV_{A_s}$  appear earlier (i.e. around 5 days) for the case of  
 610  $k_{so}=10^{-5}$  m/s compared to the case of  $k_{so}=10^{-6}$  m/s. This is because the soil with  $k_{so}=10^{-5}$   
 611 m/s is permeable enough to allow the rapid advancement of the wetting front normal to  
 612 the slope face. Instead, in the case of  $k_{so}=10^{-6}$  m/s, the narrow water path parallel to the  
 613 slope face requires a longer time to accumulate enough water at the toe slope for inducing  
 614 failure.

615

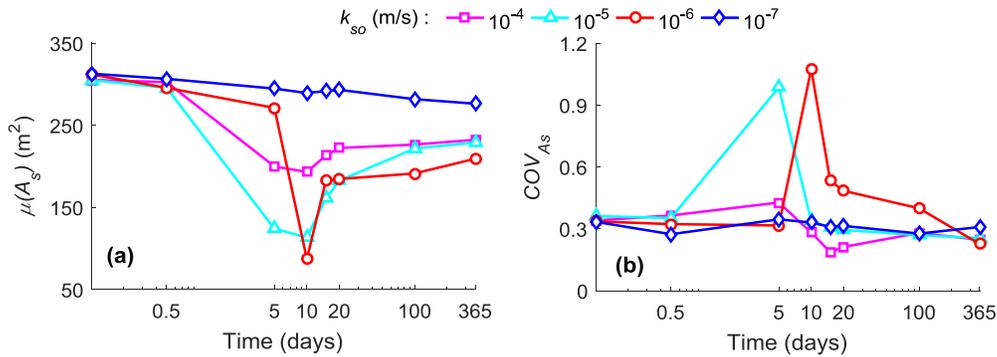


616

617 Figure 7. Time evolution of  $FoS$  in terms of mean (a) and coefficient of variation (b).

618 Analyses: influence of reference saturated permeability  $k_{so}$ .

619

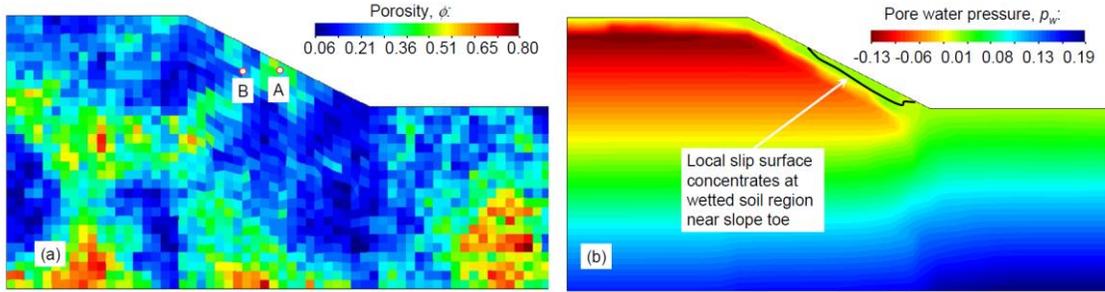


620

621 Figure 8. Time evolution of  $A_s$  in terms of mean (a) and coefficient of variation (b).

622 Analyses: influence of reference saturated permeability  $k_{so}$ .

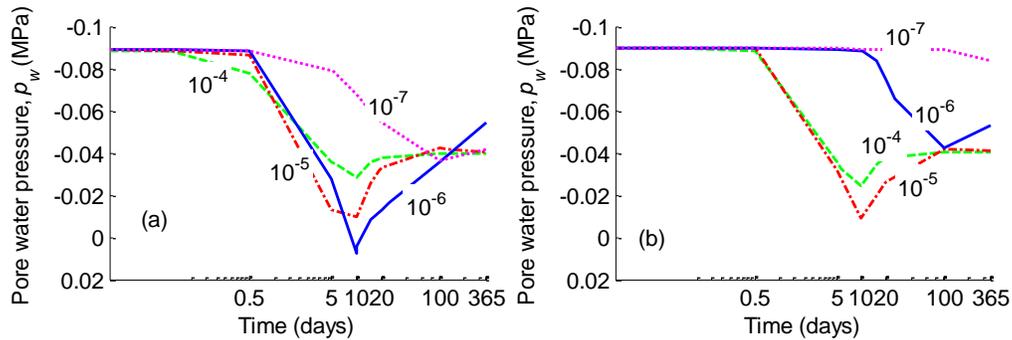
623



624  
 625 Figure 9. Porosity distribution of a sample realisation showing sampling points (a) and  
 626 contour map of  $p_w$  with slip surface at 5 days for the case of  $k_{so}=10^{-6}$  m/s (b).

627

628



629  
 630 Figure 10. Time evolution of  $p_w$  for different values of the reference saturated  
 631 permeability  $k_{so}$  at sampling points A (a) and B (b). Results correspond to the porosity  
 632 distribution and sampling points shown in Figure 9a.

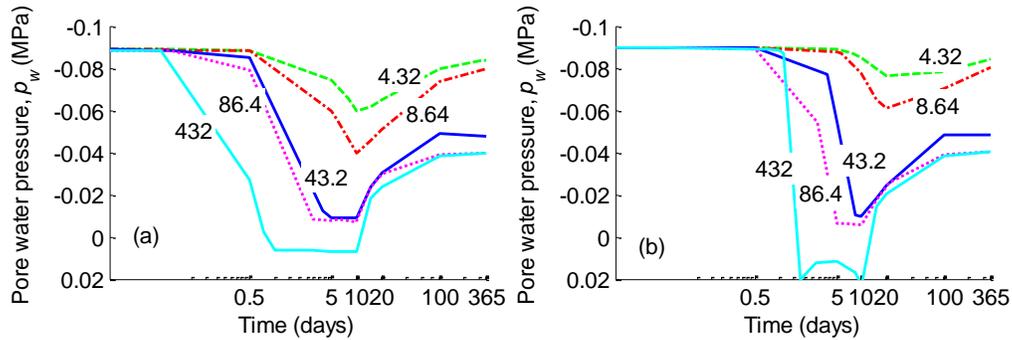
### 633 4.3 Rainfall intensity

634 The rainfall intensity  $I_r$  affects both the amount and rate of water infiltrating into the soil.  
 635 To investigate this aspect, five rainfalls of different intensities, from very light (i.e.  
 636  $I_r=4.32$  mm/day) to extremely heavy (i.e.  $I_r=432$  mm/day), are applied to each realisation  
 637 in five separate finite element simulations.

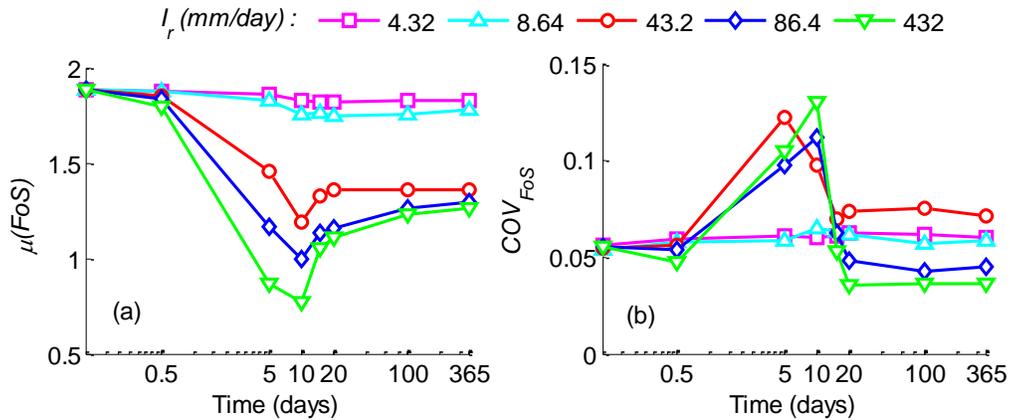
638

639 As expected, the suction drop is more significant for the heavier rainfalls as the amount  
 640 of water supply is larger (Figure 11). Therefore, the value of  $\mu(ForS)$  generally decreases  
 641 with increasing  $I_r$  with the most noticeable differences between 5 to 20 days (Figure 12a).

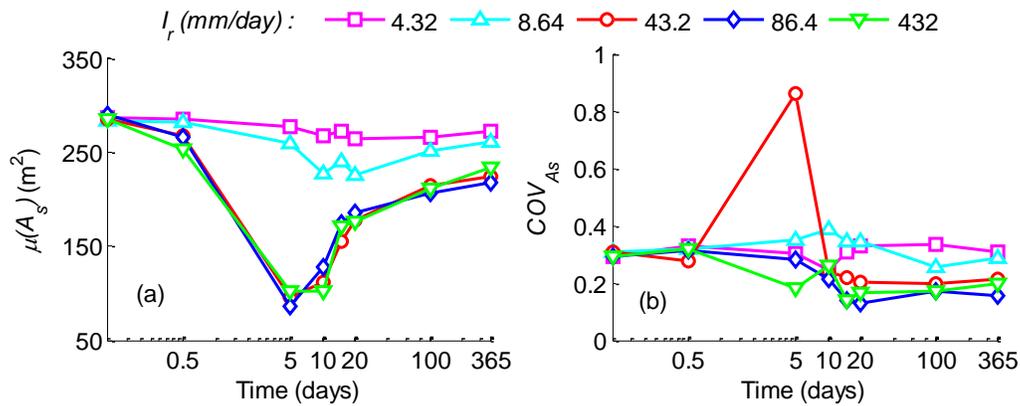
642 The two lighter rainfalls (i.e.  $I_r=4.32$  and  $8.64$  mm/day) do not provide enough water to  
 643 induce a substantial change of soil suction, hence the values of  $\mu(FoS)$ ,  $\mu(A_s)$ ,  $COV_{FoS}$  and  
 644  $COV_{A_s}$  remain approximately constant over time (Figure 12 and Figure 13).  
 645



646  
 647 Figure 11. Time evolution of  $p_w$  for different rainfall intensities  $I_r$  at sampling points A  
 648 (a) and B (b). Results correspond to the porosity distribution and sampling points shown  
 649 in Figure 9a.



650  
 651 Figure 12. Time evolution of  $FoS$  in terms of mean (a) and coefficient of variation (b).  
 652 Analyses: influence of rainfall intensity  $I_r$ .  
 653



654

655 Figure 13. Time evolution of  $A_s$  in terms of mean (a) and coefficient of variation (b).

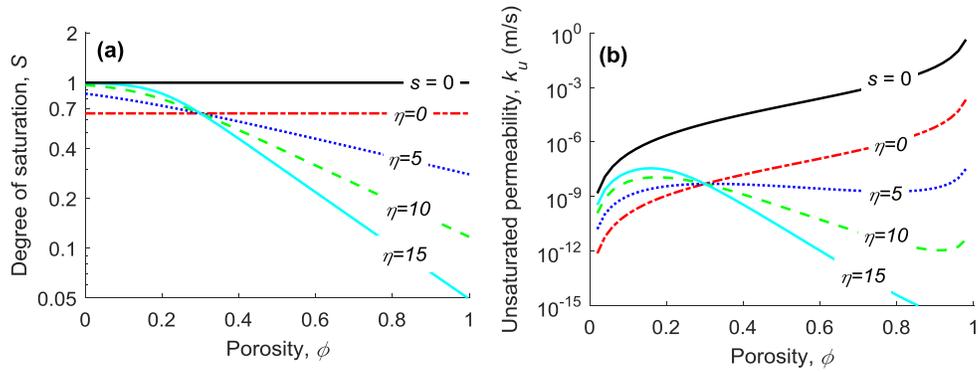
656 Analyses: influence of rainfall intensity  $I_r$ .

#### 657 4.4 Soil water retention curve – Parameter $\eta$

658 The parameter  $\eta > 0$  controls the dependency of the air entry value  $s_e$  (Eq. 2) on porosity  
 659 and therefore influences the variation of both degree of saturation  $S$  (Eqs. 1 and 2) and  
 660 unsaturated permeability  $k_u = k_r k_s$  (Eqs. 1, 2 and 4) with porosity. Figure 14 shows the  
 661 variation of degree of saturation  $S$  and unsaturated permeability  $k_u$  with porosity  $\phi$  at a  
 662 reference suction  $s = 100$  kPa for four different values of  $\eta$ , namely  $\eta = 0, 5, 10$  and  $15$ . The  
 663 non-monotonic variation of unsaturated permeability  $k_u$  (Figure 14b) is the result of the  
 664 competition between the growth of saturated permeability  $k_s$  (Eq. 3) and the reduction of  
 665 relative permeability  $k_r$  (Eq. 4) with increasing porosity  $\phi$ . For  $\eta = 0$ , however, the  
 666 variation of unsaturated permeability  $k_u$  with porosity  $\phi$  is exclusively governed by the  
 667 saturated permeability  $k_s$  as the degree of saturation  $S$ , and hence the relative permeability  
 668  $k_r$ , are independent of porosity. This explains the monotonic variation of  $k_u$  for the  
 669 particular case where  $\eta = 0$  (Figure 14b).

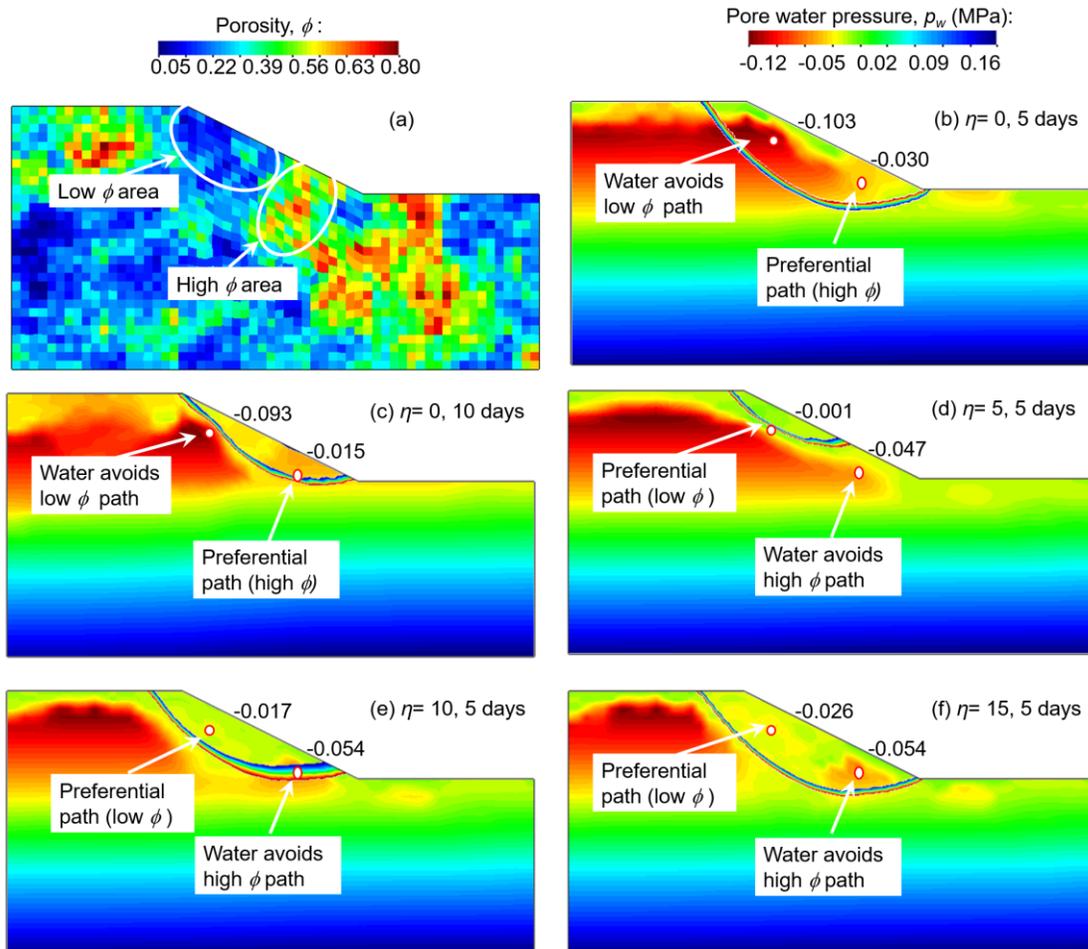
670

671 In Figure 14, the curves for different values of  $\eta$  cross each other at the reference porosity  
 672  $\phi_0$ , which means that for  $\phi > \phi_0$  the degree of saturation  $S$  and the unsaturated permeability  
 673  $k_u$  increase with increasing  $\eta$  while the opposite is true for  $\phi < \phi_0$ .



674

675 Figure 14. Variation of  $S$  (a) and  $k_u$  (b) with  $\eta$  at a reference suction  $s=100$  kPa. For the  
 676 saturated case (i.e.  $s=0$ ),  $S$  and  $k_u$  are independent of  $\eta$ .



677

678 Figure 15. Porosity distribution of a sample realisation (a) and corresponding contour  
 679 maps of  $p_w$  with slip surfaces at different times and for different  $\eta$  values (b, c, d, e, f).

680 For  $\eta=5, 10$  or  $15$ , the reduction of suction caused by rainfall infiltration is more  
681 significant in the low porosity regions (i.e. in the upper part of the slope for the  
682 realisation shown in Figure 15a) than in the high porosity ones (Figures 15d, 15e, 15f)  
683 while the opposite is true for  $\eta=0$  (Figures 15b, 15c). This is because, when  $\eta=5, 10$  or  
684  $15$ , the water preferentially flows through low porosity regions, i.e. those regions where  
685  $\phi < \phi_0$ , due to their higher unsaturated permeability (Figure 14b). The opposite is true for  
686 the case where  $\eta=0$ .

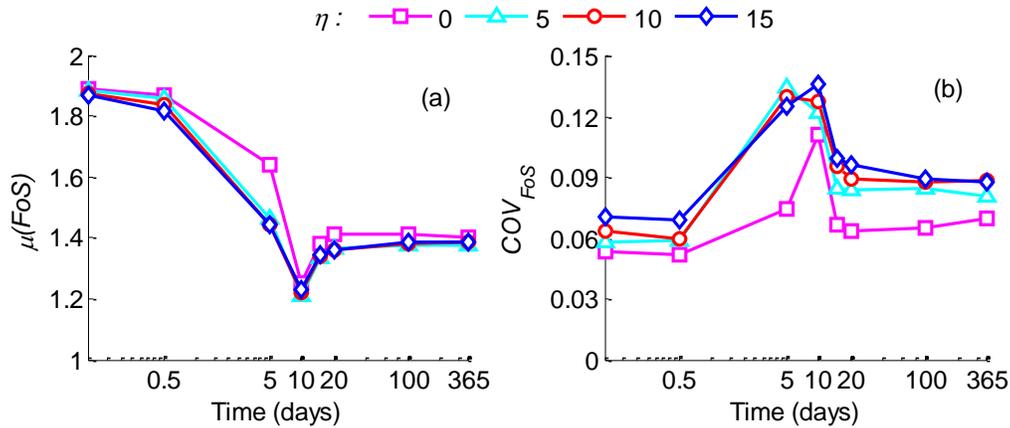
687

688 Figure 16a shows the variation of  $\mu(FoS)$  with time, which is almost identical for the  
689 three cases where  $\eta=5, 10$  or  $15$  and significantly bigger for the case where  $\eta=0$ . This  
690 pattern is justified by the fact that, in the absence of coupling between porosity and air  
691 entry value (i.e.  $\eta=0$ ), water flows preferentially through the higher porosity regions,  
692 which require longer times to become saturated. This delays the advancement of the  
693 wetting front and explains the higher values of  $\mu(FoS)$  for  $\eta=0$  compared to  $\eta=5, 10$  or  
694  $15$ . The values of  $COV_{FoS}$  are also relatively similar for the three cases where  $\eta=5, 10$  or  
695  $15$  but significantly smaller for the case where  $\eta=0$  (Figure 16b).

696

697 In terms of sliding area, the value of  $\mu(A_s)$  decreases with decreasing  $\eta$ , except for the  
698 case where  $\eta=0$ , which exhibits the highest value of  $\mu(A_s)$  at 5 days due to the delayed  
699 advancement of the wetting front (Figure 17a). The unsaturated permeability  $k_u$  exhibits  
700 the weakest dependency on porosity  $\phi$  for the case where  $\eta=5$  (Figure 14) leading to  
701 similar reductions of suction in the superficial wetted region regardless of whether  
702 porosity is high or low. This also explains why, in the case of  $\eta=5$ , suction is lower and  
703 full saturation of the top layer is reached at around 5 days (Figure 15d), leading to the  
704 formation of smaller sliding areas, i.e. lower values of  $\mu(A_s)$  and higher values of  $COV_{A_s}$   
705 (Figure 17 b).

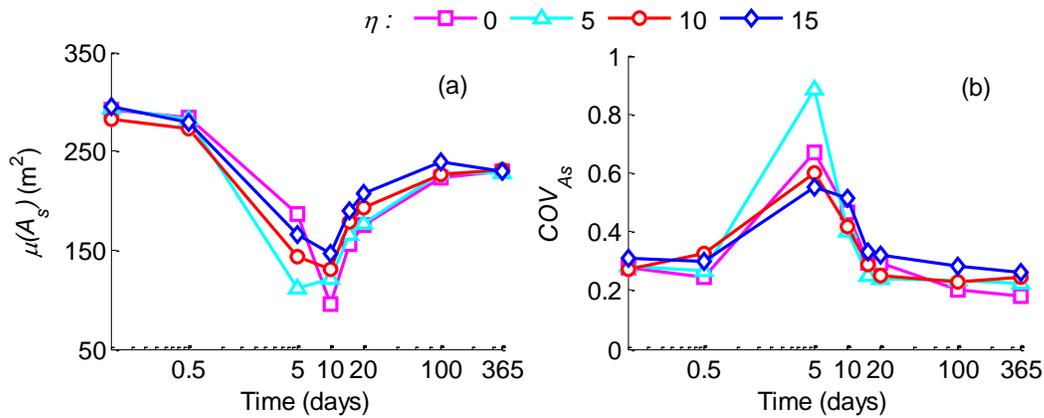
706



707

708 Figure 16. Time evolution of  $FoS$  in terms of mean (a) and coefficient of variation (b).

709 Analyses: influence of the SWRC (parameter  $\eta$ ).



710

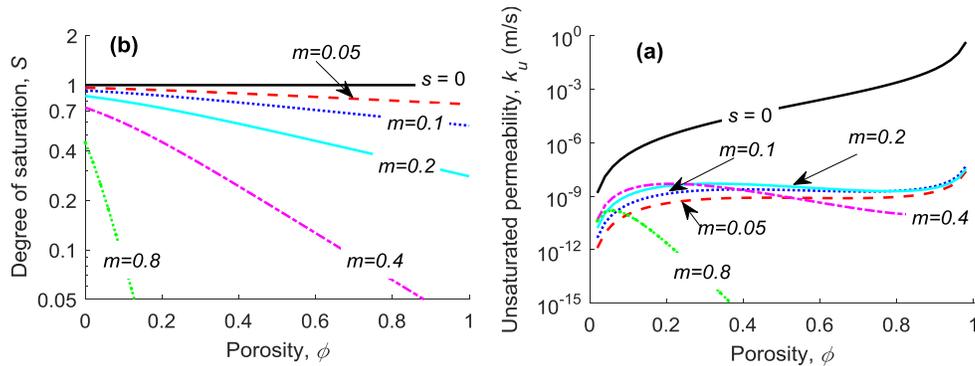
711 Figure 17. Time evolution of  $A_s$  in terms of mean (a) and coefficient of variation (b).

712 Analyses: influence of the SWRC (parameter  $\eta$ ).

### 713 4.5 Soil water retention curve – Parameter $m$

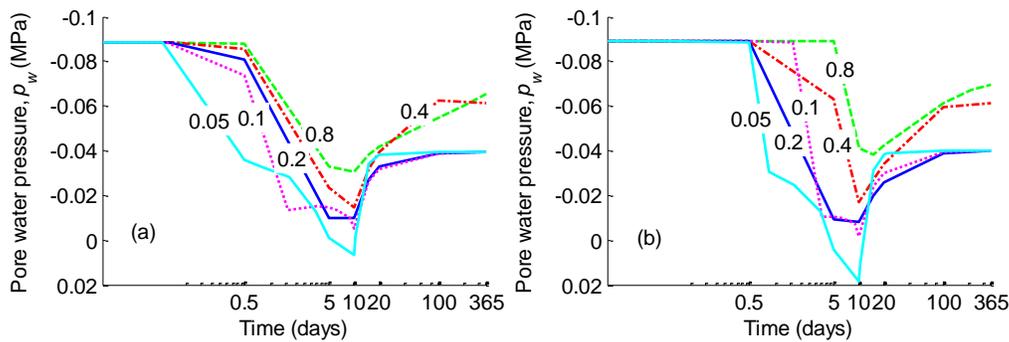
714 The slope of the water retention curve (Eq. 1) becomes more pronounced as the value of  
 715 parameter  $m$  increases, which results in a decrease of degree of saturation and unsaturated  
 716 permeability at a given suction (Eqs. 1 and 4). Figure 18 shows the variation of degree of  
 717 saturation  $S$  and unsaturated permeability  $k_u = k_r k_s$  with porosity  $\phi$  at a reference suction  
 718  $s = 100$  kPa for four different values of  $m$ , namely  $m = 0.05, 0.1, 0.2, 0.4$  and  $0.8$ . The  
 719 variation of  $k_u$  with  $\phi$  is relatively modest for  $m \leq 0.4$  because of the competing effects of

720 the increase of saturated permeability  $k_s$  (Eq. 3) and the decrease of relative permeability  
 721  $k_r$  (Eq. 4) with increasing porosity  $\phi$ .



722  
 723 Figure 18: Variation of  $S$  (a) and  $k_u$  (b) with  $m$  at a reference suction  $s=100$  kPa. For the  
 724 saturated case (i.e.  $s=0$ )  $S$  and  $k_u$  are independent of  $m$ .

725 For a given porosity, if the value of  $m$  is small, the soil exhibits a high initial value of  $S$   
 726 and therefore requires less water to reach the saturated state (Figure 18a). This produces a  
 727 quicker advancement of the wetting front so that an earlier and larger reduction of suction  
 728 occurs in the superficial soil layer as shown in Figure 19. This in turn causes an earlier a  
 729 larger reduction of shear strength, which explains why at the end of the rainfall (i.e. 10  
 730 days) the value of  $\mu(FoS)$  is about 1.6 for  $m=0.8$  but less than 1 for  $m=0.05$  (Figure 18a).



731  
 732 Figure 19. Time evolution of  $p_w$  for different values of parameter  $m$  at sampling points A  
 733 (a) and B (b). Results correspond to the porosity distribution and sampling points shown  
 734 in Figure 6b.

735 In Figure 20b, the value of  $COV_{FoS}$  increases with increasing  $m$  at initial times (i.e.  
 736 between 0 and 0.5 day) because of the increasing variability in overburden weight.  
 737 However, the highest  $COV_{FoS}$  is achieved at 5 days for an intermediate value of  $m=0.2$ ,

738 which produces the largest spread of failure mechanisms (e.g. Figure 6c and Figure 6d).  
739 This is also reflected in the relatively large value of  $COV_{A_s}$ . For the larger value  $m = 0.4$ ,  
740 the value of  $COV_{FoS}$  peaks at 10 days instead of 5 days due to the slower migration of the  
741 wetting front compared to the case of  $m = 0.2$  as discussed earlier. Similarly, the  
742 magnitude of the peak is smaller because most realisations have not reached yet the  
743 critical depth. For the smaller values  $m = 0.05$  and  $0.1$ , the wetting front advances faster  
744 and is likely to have already passed the critical depth at 5 days. At this time, the vast  
745 majority of realisations therefore exhibit sliding areas confined to the top wetted region  
746 and correspond, on average, to lower values of  $FoS$  and  $A_s$ . In this case, the peak of  
747  $COV_{FoS}$  at 10 days is caused by the development of a different failure mechanism caused  
748 by the rise of the water table in a considerable number of realisations. This higher water  
749 table produces the build-up of positive pore pressures and the formation of slip surfaces  
750 cutting through the deep saturated region.

751

752 As for the largest value  $m = 0.8$ , the  $COV_{FoS}$  uncharacteristically drops to the lowest value  
753 at 10 days (Figure 20b). This is probably due to the fact that the rainfall infiltration  
754 reduces the initially large non-uniformity of overburden weight in the unsaturated zone.

755

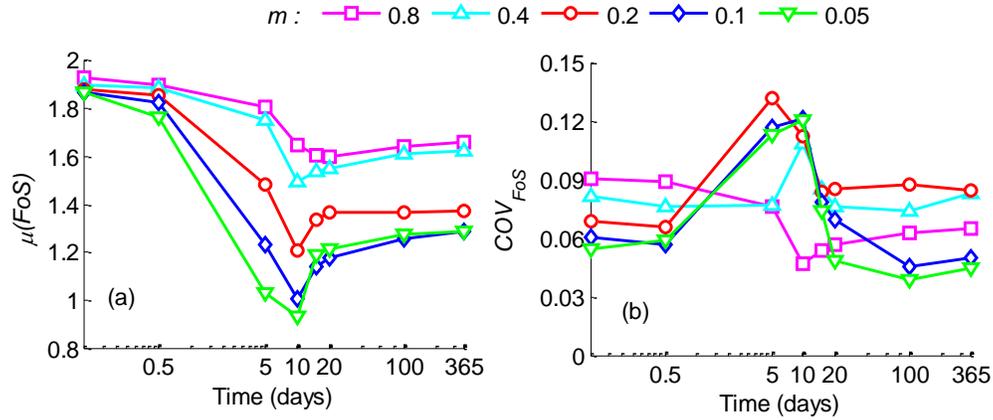
756 The value of  $\mu(A_s)$  decreases during the rainfall with the lowest values recorded between  
757 5 days for  $m = 0.1$  and 10 days for  $m = 0.05, 0.2, 0.4$  and  $0.8$  (Figure 21a). The values of  
758  $\mu(A_s)$  for  $m = 0.4$  and  $0.8$  are generally higher than in all other cases because the wetting  
759 front did not reach the critical depth in the majority of realisations, which means that the  
760 factor of safety and sliding area are generally large.

761

762 The variation of  $A_s$  between realisations is marginal for small values of  $m$  (i.e.  $0.05$  and  
763  $0.1$ ) with no prominent peaks of  $COV_{A_s}$  (Figure 21b). The fast advancement of the  
764 wetting front suggests that, in these cases, the peaks might have occurred between  $0.5$  and  
765  $5$  days, hence they are not shown in Figure 21b. Conversely, the  $COV_{A_s}$  for  $m = 0.2$   
766 exhibits a sharp peak indicating a large spread of failure mechanisms at  $5$  days and hence  
767 a large variation of  $A_s$  between realisations as previously discussed. As before, the slower

768 advancement of the wetting front delays the attainment of the peak value of  $COV_{A_s}$  to 10  
 769 days for the two cases of  $m = 0.4$  and  $0.8$  (Figure 21b).

770

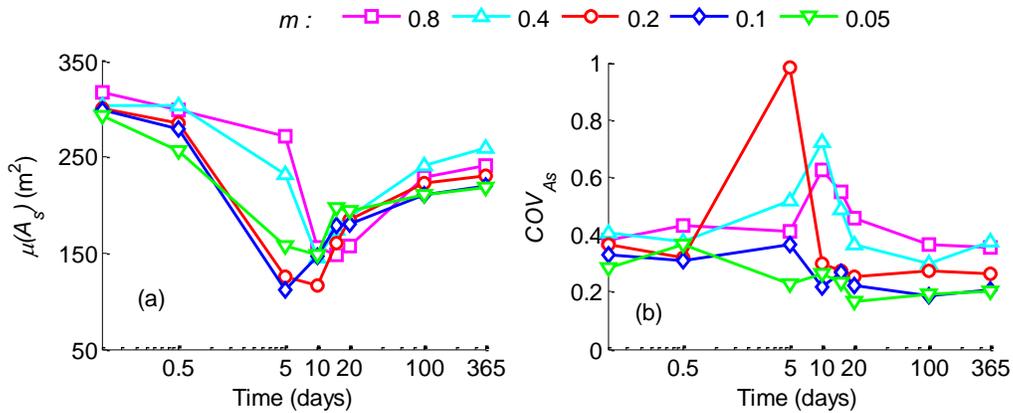


771

772 Figure 20: Time evolution of  $FoS$  in terms of mean (a) and coefficient of variation (b).

773 Analyses: influence of the SWRC (parameter  $m$ ).

774



775

776 Figure 21: Time evolution of  $A_s$  in terms of mean (a) and coefficient of variation (b).

777 Analyses: influence of the SWRC (parameter  $m$ ).

778

## 779 **5 Conclusions**

780 This study has shown that the interaction between randomly heterogeneous porosity and  
781 partial saturation can lead to very complex statistical variations of both factor of safety  
782 and failure size in soil slopes exposed to rainfall infiltration. In general, infiltration  
783 diminishes the stability of an unsaturated slope but the extent of this effect depends on  
784 various factors. If the slope exhibits large porosity variability, results can change  
785 significantly among realisations and fluctuate considerably over time, which may lead to  
786 different conclusions about the safety of the slope compared to the homogeneous case.  
787 Moreover, the statistical variation of the factor of safety and failure size is strongly  
788 influenced by other factors such as water table depth, rainfall intensity, saturated  
789 permeability and retention parameters.

790

791 The advancement of the wetting front during rainfall has a strong influence on both factor  
792 of safety and failure size. If the wetting front attains or surpass a 'critical' depth, failure is  
793 confined within the wetted superficial layer with a relatively low factor of safety.  
794 Conversely, if the wetting front is shallower than the critical depth, the failure surface  
795 penetrates deep in the soil, through both wetted and unwetted regions, with a relatively  
796 high factor of safety. During rainfall, the mean values of both factor of safety and failure  
797 size decrease because of the progressive reduction of soil suction in the superficial soil  
798 layer. These mean values attain their respective minima when the majority of Monte  
799 Carlo realisations exhibit wetting fronts deeper than the critical depth. After the end of  
800 the rainfall, these mean values increase again as suction is progressively recovered. The  
801 coefficients of variation of both factor of safety and failure size also increase until the  
802 wetting front attains the critical depth in a significant number of realisations. At this time,  
803 the failure mechanism may vary widely from shallow to deep seated, which produces  
804 large coefficients of variation.

805

806 An increase in rainfall intensity leads to a faster drop in suction, which elevates the risk  
807 of failure. Conversely, a progressive increase of saturated permeability only elevates the  
808 risk of failure up to a limit, after which the probability of failure starts to reduce. This is

809 because a very high permeability allows excess pore water pressures to dissipate quickly  
810 while a very low permeability impedes infiltration altogether. Both these effects decrease  
811 the possibility of failure, which explains why the highest risk corresponds to an  
812 intermediate permeability level.

813

814 The effect of porosity on unsaturated permeability is non-monotonic due to the opposite  
815 variation of the saturated and relative permeability. This complex behaviour produces  
816 rather unexpected patterns of water flow in heterogeneous unsaturated slopes. If the  
817 retention curve is independent of porosity, water preferably migrates through high  
818 porosity regions but, if a pronounced dependency on porosity is introduced, water tends  
819 to move through low porosity areas. Moreover, the risk of failure is significantly higher if  
820 a dependency of water retention on porosity is assumed and if the gradient of the  
821 retention curve is small to intermediate.

822

823 The progressive infiltration of water reduces both factor of safety and sliding area. This  
824 does not mean that a large sliding cannot occur in correspondence of a low factor of  
825 safety but only means that a small failure might initially occur triggering a progressively  
826 larger mechanism. It also suggests that a more accurate assessment of risk should be  
827 based on the likelihood of both slope failure and large sliding area.

828

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